

Remarks on Spectral Equivalence of Certain Discrete Operators

Włodzimierz Proskurowski*

Abstract We consider the Neumann-Dirichlet preconditioner for the discrete Laplacian in the unit square. We show that the capacitance matrix C is equal to the identity even for problems with discontinuous coefficients. In the case of many Neumann-Dirichlet strips this is no longer true if the strips are extremely thin. The conditioning of C in this case is significantly better when the Neumann strips correspond to regions with larger coefficients.

1. Introduction It is well known that all uniformly elliptic operators L defined on Ω with the same boundary conditions are *spectrally equivalent* [5]:

$$c_1(L_1x, x) \leq (L_2x, x) \leq c_2(L_1x, x),$$

where c_1, c_2 are positive constants, and (x, y) is a proper inner product. Similar relations hold for the discretized form of these operators:

$$a_1x^T Ax \leq x^T Bx \leq a_2x^T Ax, \quad \forall x \in \mathfrak{R}^n, \quad \forall n,$$

where A, B are $n \times n$ symmetric positive definite matrices, and a_1, a_2 are positive constants independent of n . These inequalities imply that the ratio of extreme eigenvalues of AB^{-1} is bounded by a_2/a_1 , called the *spectral equivalence bound* [1]. Thus, we could use one discrete elliptic operator as an efficient preconditioner for another one. In particular, the discrete Laplacian would be an excellent candidate for such a preconditioner. Let us have a closer look at some questions that can be posed: How large is the bound for $\kappa(AB^{-1})$? What is the effect of changing the boundary conditions? And finally, how all this relates to domain decomposition?

*Dept. of Mathematics, University of Southern California, Los Angeles, CA 90089-1113.

1.1 Changing boundary conditions Let us consider two $n \times n$ matrices, A and B , that represent the one dimensional -Laplacian on $\Omega = [0,1]$ with the Dirichlet and Neumann boundary condition at the left end of the interval, and with the Dirichlet condition at the other end, respectively. They differ only in the (1,1) element:

$$A = \begin{pmatrix} 2 & -1 & & & \\ -1 & 2 & -1 & & \\ & & \cdot & \cdot & \cdot \\ & & & -1 & 2 & -1 \\ & & & & -1 & 2 \end{pmatrix}, \quad B = \begin{pmatrix} 1 & -1 & & & \\ -1 & 2 & -1 & & \\ & & \cdot & \cdot & \cdot \\ & & & -1 & 2 & -1 \\ & & & & -1 & 2 \end{pmatrix}.$$

It has been shown by Hald, see [7, p.457], that AB^{-1} is a rank one modification of an identity, $AB^{-1} = I + uv^T$, where $u^T = (1,0,\dots,0)$ and $v^T = (n,n-1,\dots,1)$. Moreover, only two singular values of AB^{-1} differ from 1. These two coincide with the eigenvalues of $I_2 + (v+\alpha u, u)^T(u, v + \alpha u)$, where $\alpha = v^T v / 2 \approx n^3/6$. Its characteristic polynomial is $\lambda^2 - 2(1+\alpha+\gamma)\lambda + (1+\gamma^2+2\gamma) = 0$, where $\gamma = u^T v = n$. Thus, the smallest singular value of AB^{-1} is equal to about $\sigma_{\min} \approx 3/n$, the largest $\sigma_{\max} \approx n^3/3$, and $\kappa(AB^{-1}) = \sqrt{\sigma_{\max} / \sigma_{\min}} \approx n^2/3$. Consequently, A and B are far from being spectrally equivalent, although they represent the same operator (-Laplacian), albeit with different boundary conditions.

2. Neumann-Dirichlet preconditioner (two subdomains) Let us now consider the same -Laplacian on $\Omega = [0,1]$ with the Dirichlet boundary conditions at both end points, represented by the $n \times n$ matrix A . Let us impose artificial inner boundary conditions at $x=0.5$ such that we have the Dirichlet condition to the left of it, and the Neumann condition to the right. The resulting matrix B has the form:

$$B = \begin{pmatrix} 2 & -1 & & & \\ \cdot & \cdot & \cdot & & \\ & -1 & 2 & -1 & \\ & & 0 & 2 & -2 \\ & & & -1 & 2 & -1 \\ & & & & \cdot & \cdot & \cdot \\ & & & & & -1 & 2 \end{pmatrix}$$

Using the same technique as above, we can show that $AB^{-1} = I + uv^T$, where $u^T = (0,\dots,0,1,0,\dots,0)$, i.e., $u_i = \delta_{ik}$, $v^T = 1/k (-1,\dots,-(k-1),0,k-1,\dots,1)$, and $k = (n+1)/2$. Since now $\gamma = u^T v = 0$, the characteristic polynomial of $I_2 + (v+\alpha u, u)^T(u, v + \alpha u)$ is simplified to $\lambda^2 - 2(1+\alpha)\lambda + 1 = 0$. Here, the new $\alpha = v^T v / 2 \approx n/3$. Thus, $\sigma_{\min} \approx 3/n$, similarly as before, but $\sigma_{\max} \approx n/3$. As a result, we obtain a much more favorable ratio,

$N=n^2$	λ_{\min}	λ_{\max}	cond #
9	.76	1.31	1.72
25	.66	1.51	2.28
49	.60	1.68	2.82
81	.55	1.83	3.35

Table 2. Extreme eigenvalues and the condition number of $\mathbb{A}\mathbb{B}^{-1}$ in 2D.

Table 2 demonstrates the results of numerical experiments. Here, the condition number grows linearly with \sqrt{N} (approximately as $\sqrt{N/4+1}$), yet, once more, \mathbb{A} and \mathbb{B} fail to be spectrally equivalent.

It should be noted, that application of a special symmetrizer to \mathbb{B} , see [8], results in a matrix \mathbb{B} that is spectrally equivalent to \mathbb{A} . More precisely, only n eigenvalues of $\mathbb{A}\mathbb{B}^{-1}$ are equal to 2, while all the rest are 1.

3. Capacitance matrix The method of domain decomposition often can be considered as a process in a *subspace*, see [6]. This amounts to performing the main iteration with the capacitance matrix C of the form $C = S^T \mathbb{A} \mathbb{B}^{-1} S$, where S^T is a restriction operator $S^T = (I_p, 0)$, and p is the number of grid points on the separator, $p \ll n$. Note, that for our one dimensional examples in sections 1.1 and 2., p equals to 1 and the capacitance $C = S^T (I + uv^T) S$ is equal to $n+1$ and 1 , respectively.

In general, we can write matrices \mathbb{A} and \mathbb{B} in a 2 by 2 block form:

$$\mathbb{A} = \begin{pmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{pmatrix} \quad \mathbb{B} = \begin{pmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{pmatrix} \quad (2)$$

It has been shown, see [3] and also [4], that if A_{11} is invertible then $\mathbb{A}\mathbb{B}^{-1}$ has the form:

$$\mathbb{A}\mathbb{B}^{-1} = \begin{pmatrix} I & 0 \\ Z & C \end{pmatrix} \quad (3)$$

where $C = C_1 C_2^{-1}$, $C_1 \equiv (S^T \mathbb{A} S)^{-1} = A_{22} - A_{21} A_{11}^{-1} A_{12}$,

$C_2 \equiv (S^T \mathbb{B} S)^{-1} = B_{22} - B_{21} A_{11}^{-1} A_{12}$, and $Z = (A_{21} - C B_{21}) A_{11}^{-1}$.

C_1 and C_2 are called Schur complements of \mathbb{A} and \mathbb{B} , respectively, and thus it is appropriate to call C the Schur complement of $\mathbb{A}\mathbb{B}^{-1}$.

Matrix \mathbb{B} defined by (1'') can be used efficiently as the so-called Neumann-Dirichlet preconditioner for \mathbb{A} , see [2]. We want to show not only that \mathbb{A} and \mathbb{B} are spectrally equivalent in a *subspace* :

$a_1 x^T S^T \mathbb{A} S x \leq x^T S^T \mathbb{B} S x \leq a_2 x^T S^T \mathbb{A} S x, \quad \forall x \in \mathfrak{R}^n, \quad \forall n$

but also find the proportionality constants a_1, a_2 .

When we reorder \mathbb{A} and \mathbb{B} from (1) into the block form (2) we obtain:

$$\mathbb{A} = \begin{pmatrix} T & -I & & & & & & & & & 0 \\ -I & T & -I & & & & & & & & 0 \\ & & \cdot & \cdot & \cdot & & & & & & \cdot \\ & & & -I & T & -I & & & & & 0 \\ & & & & -I & T & & & & & -I \\ & & & & & T & -I & & & & -I \\ & & & & & -I & T & -I & & & 0 \\ & & & & & & \cdot & \cdot & \cdot & & \cdot \\ & & & & & & & -I & T & -I & 0 \\ & & & & & & & & -I & T & 0 \\ 0 & 0 & 0 & 0 & -I & -I & 0 & 0 & 0 & 0 & T \end{pmatrix}$$

where matrix A_{11} is $(n^2-n) \times (n^2-n)$, $A_{22}=T$ is $n \times n$, and $A_{12}=A_{21}^T$ is $(n^2n) \times n$.

Matrix \mathbb{B} differs from \mathbb{A} only in the (21) block as $B_{22} = A_{22} = T$:

$$B_{21} = (0 \quad 0 \quad 0 \quad 0 \quad 0 \quad -2I \quad 0 \quad 0 \quad 0 \quad 0 \quad 0).$$

A_{11} is a 2×2 block diagonal matrix with two identical diagonal blocks. As a consequence, so is its inverse. Let us denote $A_{11}^{-1} = \text{diag}(G, G)$, where the $m \times m$ matrix G represents the discrete Green's function for Δ_n on a half of the unit square with Dirichlet boundary conditions. The action of A_{21} on $A_{11}^{-1} A_{12}$ is identical to that of B_{21} on $A_{11}^{-1} A_{12}$ and is equal to $2S^T G S$, where the $m \times n$ matrix S is defined by $S^T = (0, \dots, 0, I)^T$. Therefore, $C_1 = C_2 = T - 2S^T G S$, and finally we get $C = C_1 C_2^{-1} = I$.

We have thus shown that C , the Schur complement of AB^{-1} for the case of the Neumann-Dirichlet preconditioner, is equal to the identity. Consequently, in a subspace, \mathbb{B} is an excellent preconditioner for \mathbb{A} .

Additional item. Matrix $Z = (A_{21} - CB_{21})A_{11}^{-1} = (0, \dots, 0, -I, I, 0, \dots, 0) \text{diag}(G, G) = (X, -X)$, where the $n \times m$ matrix $X = -S^T G$. Thus, $\mathbb{A} \mathbb{B}^{-1} = I + SZ^T$.

4. Many Neumann-Dirichlet strips In the case of many Neumann-Dirichlet strips the situation is somewhat different: the capacitance C is not equal to the unity matrix any more. Actually, it is not even symmetric. Nevertheless, since C_1 and C_2 remain symmetric, all eigenvalues of C are real, as the following argument shows.

$$C\phi = C_1 C_2^{-1} \phi = \lambda \phi, \quad C_2^{-1/2} C_1 C_2^{-1/2} \phi = \lambda \phi, \quad \text{where } C_2^{-1/2} \phi = \phi.$$

Table 3 demonstrates the dependence of the number of grid points across and along each strip on the condition number of C in the case of four strips. For very narrow strips, $\kappa(C)$ is not independent of the grid size h . In the extreme case of only one inner grid point across the strip, $\kappa(C)$ grows roughly linearly with the number of grids in the other direction, i.e., with the inverse of the grid size, $1/h$. On the other hand, when the number of grid points across the strip grows even slightly, $\kappa(C)$ rapidly decays to one.

	k=1	k=2	k=3	k=4	k=5	k=7	k=9	k=11
m=1	1.226	1.520	1.926	2.431	3.030	4.510	6.375	8.631
m=2	1.056	1.171	1.353	1.590	1.874	2.570	-	-
m=3	1.015	1.062	1.153	1.283	1.446	-	-	-
m=4	1.004	1.023	1.070	1.145	-	-	-	-

Table 3. Condition number of the capacitance matrix in the case of four strips, $d=4$. Here, k and m are the number of inner grid points along and across the strip, respectively.

In the simplest example with four strips, when $m=k=1$, we have

$$C_1 = 1/4 \begin{pmatrix} 14 & -1 & 0 \\ -1 & 14 & -1 \\ 0 & -1 & 14 \end{pmatrix} \quad C_2 = 1/4 \begin{pmatrix} 14 & -2 & 0 \\ -2 & 14 & 0 \\ 0 & 0 & 14 \end{pmatrix}$$

$$C = C_1 C_2^{-1} = 1/96 \begin{pmatrix} 97 & 7 & 0 \\ 7 & 97 & -48/7 \\ -1 & -7 & 96 \end{pmatrix}$$

with the eigenvalues equal to 1.1130, 1.0000 and 0.9078, which results in $\kappa(C)=1.226$.

We note that out of total $k(d-1)$ eigenvalues of C only k of them are exactly equal to 1, the rest are unequally clustered around the value of 1. For example, for $k=5$ and $m=3$ we have $\lambda_{\max}=1.22$, $\lambda_{\min}=0.85$, and the remaining 13 eigenvalues are between 0.97 and 1.03.

	d=8					d=16	
	k=1	k=2	k=3	k=4	k=5	k=1	k=2
m=1	1.305	1.723	2.331	3.117	4.076	1.326	1.780
m=2	1.074	1.229	1.483	-	-	1.078	-
m=3	1.019	1.082	-	-	-	-	-
m=4	1.005	1.031	-	-	-	-	-

Table 4. Condition number of the capacitance matrix in the case of $d=8$ and $d=16$ strips. Here, k and m are the number of inner grid points along and across the strip, respectively.

The matrix \mathbb{A} that arises from this discretization has a form similar to that in section 3.:

$$\mathbb{A} = \begin{pmatrix} T & -I & & & & & & & & & & 0 \\ -I & T & -I & & & & & & & & & 0 \\ & & \cdot & \cdot & \cdot & & & & & & & \cdot \\ & & & -I & T & -I & & & & & & 0 \\ & & & & -I & T & & & & & & -I \\ & & & & & & cT & -cI & & & & -cI \\ & & & & & & -cI & cT & -cI & & & 0 \\ & & & & & & & \cdot & \cdot & \cdot & & \cdot \\ & & & & & & & & -cI & cT & -cI & 0 \\ & & & & & & & & & -cI & cT & 0 \\ 0 & 0 & 0 & 0 & -I & -cI & 0 & 0 & 0 & 0 & 0 & sT \end{pmatrix},$$

where the constant $c=k_2(x)/k_1(x)$ and $s=(c+1)/2$.

As before, matrix \mathbb{B} differs from \mathbb{A} only in the (21) block :

$$B_{21} = (0 \quad 0 \quad 0 \quad 0 \quad 0 \quad -2sI \quad 0 \quad 0 \quad 0 \quad 0).$$

Theorem Let $-\text{div}(k(x)\text{gradu})=f$ in $\Omega=(0,1)^2$ with Dirichlet boundary conditions at $\partial\Omega$. Let the diffusion function $k(x)$ be discontinuous at the interface across that halves Ω , and constant in each of the subdomains. Let the Neumann-Dirichlet conditions be imposed at this interface. Then the capacitance matrix C is equal to the identity.

Proof We have $A_{11}^{-1} = \text{diag}(G, 1/cG)$, where the $m \times m$ matrix G represents, as before, the discrete Green's function for Δ_n on the half of the unit square with Dirichlet boundary conditions. Here, as before, S is defined by $S^T = (0, \dots, 0, 1)^T$. Then

$$A_{11}^{-1} A_{12} = \begin{pmatrix} G & 0 \\ 0 & 1/cG \end{pmatrix} \begin{pmatrix} S \\ cS \end{pmatrix} = \begin{pmatrix} GS \\ GS \end{pmatrix}$$

$$A_{21} A_{11}^{-1} A_{12} = (S^T, cS^T) \begin{pmatrix} GS \\ GS \end{pmatrix} = (c+1)S^T GS.$$

$$B_{21} A_{11}^{-1} A_{12} = (0, (c+1)S^T) \begin{pmatrix} GS \\ GS \end{pmatrix} = (c+1)S^T GS.$$

Thus the action of A_{21} on $A_{11}^{-1}A_{12}$ is identical to that of B_{21} on $A_{11}^{-1}A_{12}$ and is equal to $(c+1)S^T GS$. Therefore, $C_1 = C_2 = sT - 2S^T GS$, and finally we get $C = C_1 C_2^{-1} = I$. Thus, C is equal to the identity, even in the case of discontinuous coefficients.

In the case of many Neumann-Dirichlet strips the results are anti-intuitive: the conditioning of the capacitance matrix C is significantly better for problems with discontinuous coefficients than with continuous ones, as the following table indicates.

	k=1	k=2	k=3	k=4	k=5	k=7	k=9	k=11
m=1	1.038	1.081	1.134	1.194	1.261	1.420	1.612	1.842
m=2	1.010	1.029	1.057	1.090	1.127	1.210	-	-
m=3	1.003	1.011	1.026	1.047	1.071	-	-	-
m=4	1.001	1.004	1.012	1.025	-	-	-	-

Table 6. Condition number of the capacitance matrix for the problem with discontinuous coefficients ($c=k_2/k_1=10$) in the case of four strips, $d=4$. Here, k is the number of inner grid points along the strip, and m is the number of grid points across the strip.

To gain some insight, let us examine the simplest example with four strips, when $m=k=1$.

$$B = \left(\begin{array}{ccc|ccc} 4 & & & -1 & & \\ & 40 & & -10 & -10 & \\ & & 4 & & -1 & \\ \hline & & & 40 & & -10 \\ -11 & & & & 22 & \\ -11 & & & & & 22 \\ & & -11 & & & 22 \end{array} \right)$$

while A is symmetric ($A_{21}=A_{12}^T$). In accordance with (3), we obtain

$$C_1 = 1/4 \begin{pmatrix} 77 & -10 & 0 \\ -10 & 77 & -1 \\ 0 & -1 & 77 \end{pmatrix}, \quad C_2 = 1/4 \begin{pmatrix} 77 & -11 & 0 \\ -11 & 77 & 0 \\ 0 & 0 & 77 \end{pmatrix},$$

$$C = C_1 C_2^{-1} = 1/528 \begin{pmatrix} 529 & 7 & 0 \\ 7 & 529 & -48/7 \\ -1 & -7 & 528 \end{pmatrix}$$

with the $\lambda_i(C)$ equal to 1.0205, 1.0000 and 0.9832, which results in $\kappa(C)=1.038$. Note that matrix C_2 is the same as the one for continuous coefficients, scaled by $(c+1)/2=11/2$.

Increasing the number of strips above 4 does not change much the conditioning of C :

	k=1	k=2	k=3	k=4	k=5	k=1	k=2
m=1	1.050	1.108	1.182	1.271	1.374	1.053	1.115
m=2	1.013	1.038	1.076	-	-	-	-

Table 7. Condition number of the capacitance matrix for the problem with discontinuous coefficients ($c=k_2/k_1=10$) in the case of $d=8$ strips. Here, k and m are the number of inner grid points along and across the strip, respectively.

Let us now investigate the influence of the ratio of the diffusion coefficients on the condition number of matrix C.

$c=k_2/k_1$	0.001	0.01	0.1	0.3	1.0	3.	10.	100.	1000.
$\kappa(C)$ (m=k=1)	1.502	1.497	1.448	1.368	1.226	1.107	1.038	1.004	1.0004
$\kappa(C)$ (m=k=3)	1.329	1.326	1.295	1.245	1.153	1.074	1.026	1.003	1.0003

Table 8. Condition number of the capacitance matrix for the problem with discontinuous coefficients k_2 and k_1 in the case of $d=4$ strips. Here, k and m are the number of inner grid points along and across the strip, respectively.

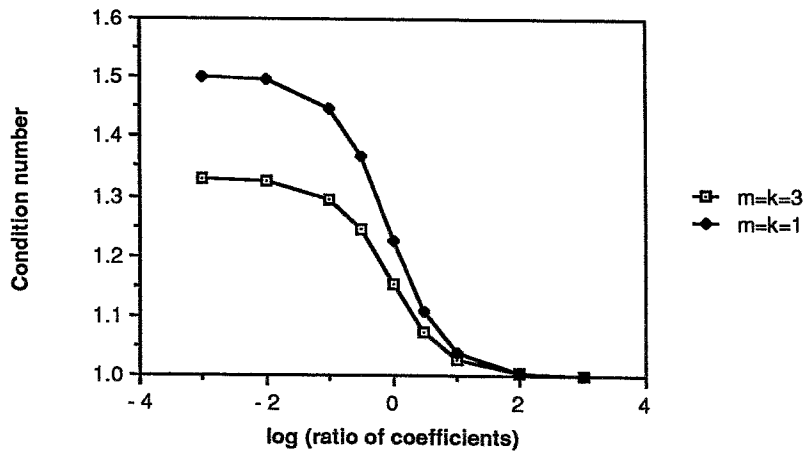
Results in Table 8 strongly indicate that as c , the ratio of the coefficients, grows, the condition number approaches the value of 1.0, and $\kappa(C)$ approaches a somewhat larger value as c decreases, see Fig.1. Coefficient k_1 corresponds here to the Dirichlet and k_2 to the Neumann strips. Let us investigate the problem for the simplest case when $k=m=1$.

Using the formulas in (3) we arrive at the following representation

$$C_1 = \frac{1}{4} \begin{pmatrix} 7(c+1) & -c & 0 \\ -c & 7(c+1) & -1 \\ 0 & -1 & 7(c+1) \end{pmatrix}, \quad C_2 = \frac{1}{4} \begin{pmatrix} 7(c+1) - (c+1) & 0 \\ -(c+1)7(c+1) & 0 \\ 0 & -1 & 7(c+1) \end{pmatrix},$$

$$C = C_1 C_2^{-1} = I + 1/48(c+1)R, \quad \text{where } R = \begin{pmatrix} 1 & 7 & 0 \\ 7 & 1 & -48/7 \\ -1 & -7 & 0 \end{pmatrix}$$

Fig.1



Thus the eigenvalues of C are $\lambda(C) = 1 + 1/48(c+1) \lambda(R)$. It is easy to verify that the eigenvalues of R are 0 (R is singular), and $1 \pm \sqrt{97}$. Consequently, the condition number of C is $\kappa(C) = (48(c+1) + 1 + \sqrt{97}) / (48(c+1) + 1 - \sqrt{97})$. This formula gives us $\kappa(C)=1.226$ for $c=1$, and $\kappa(C)=1.503$ for $c=0$ (the limiting case).

Examining C_1 and C_2 we see that for large values of c these matrices are "similar", in contrast to the situation when c are small (in the limit, for $c=0$, the zeros appear in the "wrong" positions). This strongly suggests that *in the preconditioner B the Neumann strips should correspond to the regions with larger diffusion coefficients*.

Finally, we investigate the influence of the number of inner grid points along the strip on the condition number when the strips are *wrongly* placed, i.e., $c < 1$.

k	1	2	3	4	5	7	9	11
$\kappa(C)$	1.448	2.137	3.273	4.950	7.278	14.33	25.21	40.44

Table 9. Condition number of the capacitance matrix for the problem with discontinuous coefficients ($c=k_2/k_1=0.1$) in the case of $d=4$ strips and $m=1$. Here, the k and m are the number of inner grid points along and across the strip, respectively.

Clearly, for $c < 1$, in this worst case of extremely thin strips ($m=1$), $\kappa(C)$ grows almost quadratically with k , the number of inner grid points along the strip.

References

- [1] O.Axelsson and V.A.Barker, *Finite Element Solution of Boundary Value Problems*, Academic Press, 1984.
- [2] P.Bjørstad and O.Widlund, *SIAM J. Num.Anal.*, v.23, 1987, 1097.
- [3] M.Dryja, *SIAM J.Num.Anal.*, v.20, 1983, 671.
- [4] M.Dryja and W.Proskurowski, CRI-86-12 Report, May 1986.
- [5] V.Faber, T.Manteuffel and S.Parter, "On the equivalence of operators and the implications to preconditioned iterative methods for elliptic problems", 1987, preprint.
- [6] G.Marchuk and Yu.Kuznetsov, in *Problems of Computational Mathematics and Mathematical Modelling*, G.Marchuk and V.Dymnikov (eds.), Mir Publ., 1985.
- [7] W.Proskurowski and O.Widlund, *Comp. Math.*, v.30, 1976, 433.
- [8] O.Widlund, in *Proc. 1st Symposium on Domain Decomposition*, SIAM, 1988, 113.