

Domain Decomposition Capabilities for the Mortar Finite Volume Element Methods

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Introduction

Since the introduction of the *mortar method* as a coupling technique between the spectral and finite element methods (see, e.g. [BM97, BMP94, BMSL89]), it has become the most important technique in domain decomposition methods for non-matching grids. The active research by the scientific computation community in this field is motivated by its flexibility and great potential for large scale parallel computation (see, e.g. [BM94]). A good description of the mortar element method can be found in [Bel97, BDM90, BMP94, Cas]. The nonconforming finite element mortar method has been studied in [BMP94], where optimal order convergence in H^1 -norm was demonstrated. Three-dimensional mortar finite element analysis has been given in [BM97]. Non-mortar mixed finite element approximations for second order elliptic

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problems have been discussed in [AY97].

The above mentioned mortar elements are defined on non-matching grids with non-overlapping subdomains. Recently, the overlapping mortar linear finite element method was studied in [CDS99], where several additive Schwarz preconditioners have been proposed and analyzed and extensive numerical examples to support the theoretical results have been reported.

To the authors' best knowledge, there has not been a study for the mortar finite volume element method. In the past 10 years the finite (control) volume method has drawn serious attention both from mathematicians, engineers, and physicists as an attractive solution technique for various applied problems (see, e.g. [BV96]). Following the notations and the approach of Ben Belgacem [Bel97], we apply the mortar technique to derive two control volume schemes based on: (1) finite volume element approximation of the solution in the subdomains and finite element approximation for Lagrange multipliers on the interfaces; (2) finite volume approximations for both the solution in the subdomains and the Lagrange multipliers on the interfaces. It has been shown on various test problems that the latter schemes converge much faster (5-6 times) than the former schemes. In this respect we have found evidence from our numerical experiments and we believe that in the finite volume element methods if r -th order piecewise polynomials are used on the subdomains, then $(r - 1)$ -th order polynomials should be used on the interface for the Lagrange multipliers. For both types of schemes, we have obtained in [ELLL98] optimal order H^1 -norm error estimates under the regularity assumption that $u \in H^{1+\tau_k}(\Omega_k)$ for $0 < \tau_k \leq 1$ where $\overline{\Omega} = \cup \overline{\Omega}_k$.

Mortar Finite Element Approximation

We shall use the notations from [Bel97]. We break up the initial domain Ω into K non-overlapping subdomains $\{\Omega_k\}_{1 \leq k \leq K}$, which are assumed to be polygonally shaped and arranged in such a way that the intersection of two subdomains $\overline{\Omega}_l \cap \overline{\Omega}_k$ as well as the intersection $\partial\Omega \cap \partial\Omega_k$ is either empty or reduced to a vertex or to a common edge. If two subdomains Ω_k and Ω_l are adjacent, Γ_{kl} is the common interface, and \mathbf{n}_{kl} is the unit normal from Ω_k to Ω_l . Let \bar{k} denote the set of all indices so that kl is meaningful. For any k , let $H_*^1(\Omega_k)$ denote the space $H^1(\Omega_k)$ if the measure of $\partial\Omega_k \cap \partial\Omega$ is zero; otherwise it coincides with the subspace of $H^1(\Omega_k)$ involving all functions whose trace is zero over the set $\partial\Omega_k \cap \partial\Omega$:

$$H_*^1(\Omega_k) = \{v_k \in H^1(\Omega_k) : v_k|_{\partial\Omega_k \cap \partial\Omega} = 0, \text{ if } \text{meas}(\partial\Omega_k \cap \partial\Omega) \neq 0\}.$$

Set the space

$$X = \{v \in L^2(\Omega) : v_k = v|_{\Omega_k} \in H_*^1(\Omega_k)\} = \prod_{k=1}^K H_*^1(\Omega_k)$$

equipped with the norm: $\|u\|_X = \left(\sum_{k=1}^K \|v_k\|_{H^1(\Omega_k)}^2\right)^{1/2}$. Let

$$H_0(\text{div}, \Omega) = \{q \in H(\text{div}, \Omega) : \mathbf{q} \cdot \mathbf{n}|_{\partial\Omega} = 0\},$$

where $H(\operatorname{div}, \Omega)$ is the space of all vector-functions in $(L^2(\Omega))^2$ whose weak divergence is in $L^2(\Omega)$. The trace of these function on the boundary $\partial\Omega$ is understood in the appropriate weak sense. The characterization of $H_0^1(\Omega)$ can be made:

$$H_0^1(\Omega) = \left\{ v \in X : \sum_{k=1}^K (\mathbf{q} \cdot \mathbf{n}_k, v)_{\partial\Omega_k} = 0, \quad \mathbf{q} \in H_0(\operatorname{div}, \Omega) \right\}.$$

Now we define the space M of those $\psi = (\psi_1, \dots, \psi_K)$ with components $\psi_k \in H_*^{-1/2}(\partial\Omega_k)$ such that there weak traces on the boundaries represent a weak trace of a function in $H_0(\operatorname{div}, \Omega)$, i.e.

$$M = \{ \psi : \text{there exists } \mathbf{q} \in H_0(\operatorname{div}, \Omega) \text{ s.t. for } k = 1, \dots, K, \psi_k = \mathbf{q} \cdot \mathbf{n}_k \}.$$

The space M is provided with the norm

$$\|\psi\|_M = \inf \left\{ \|\mathbf{q}\|_{H(\operatorname{div}, \Omega)} : \mathbf{q} \in H_0(\operatorname{div}, \Omega), \quad \mathbf{q} \cdot \mathbf{n}_k = \psi_k, \forall k \right\},$$

where $H_*^{-1/2}(\partial\Omega_k)$ is the dual space of $H_*^{1/2}(\partial\Omega_k)$ with $\langle \cdot, \cdot \rangle_{*, \partial\Omega_k}$ pairing, $H_*^{1/2}(\partial\Omega_k) = H^{1/2}(\partial\Omega_k)$ if $\partial\Omega_k \cap \partial\Omega = \emptyset$ and $H_*^{1/2}(\partial\Omega_k) = H_{00}^{1/2}(\partial\Omega_k \setminus \partial\Omega)$ if $\partial\Omega_k \cap \partial\Omega \neq \emptyset$. Basically speaking the constraints on the distributions $\psi \in M$ imply that the jumps across the interfaces Γ_{kl} vanish.

We now define the bilinear form: $B : X \times M \rightarrow R$ by

$$B(v, \phi) = \sum_{k=1}^K \langle v_k, \phi_k \rangle_{*, \partial\Omega_k},$$

so that it follows from Hahn-Banach Theorem that

$$H_0^1(\Omega) = \{ v \in X, \quad B(v, \phi) = 0, \quad \phi \in M \}.$$

Similarly, the bilinear form $A : X \times X \rightarrow \mathcal{R}$ is defined by

$$A(u, v) = \sum_{k=1}^K \int_{\Omega_k} \nabla u_k \cdot \nabla v_k dx.$$

We consider the following model problem: find $u \in H_0^1(\Omega)$ such that

$$A(u, v) = (f, v), \quad v \in H_0^1(\Omega). \quad (1)$$

Its primal hybrid formulation is therefore defined by: find $(u, \psi) \in X \times M$ such that

$$\begin{aligned} A(u, v) + B(v, \psi) &= (f, v), & v \in X, \\ B(u, \phi) &= 0, & \phi \in M. \end{aligned} \quad (2)$$

We have the following equivalent result: Problem (2) has a unique solution $(u, \psi) \in X \times M$, and the first component $u \in H_0^1(\Omega)$ is also the solution of problem (1). Moreover, we have

$$\psi_k = A \nabla u_k \cdot \mathbf{n}_k, \quad k = 1, \dots, K \quad \text{and} \quad \|u\|_{H_0^1(\Omega)} + \|\psi\|_M \leq C \|f\|_{L^2(\Omega)}.$$

Finite Volume Element Approximation

Let the triangulation \mathcal{T}_{h_k} of each subdomain Ω_k , $1 \leq k \leq K$, be such that

$$\bar{\Omega}_k = \cup_{T \in \mathcal{T}_{h_k}} \bar{T}, \quad h_k = \max_{T \in \mathcal{T}_{h_k}} h_T, \quad \text{and} \quad h_T = \sup_{x, y \in T} d(x, y).$$

For piecewise linear finite element subspaces of $H_*^1(\Omega_k)$ on \mathcal{T}_{h_k} , we set

$$X_{\delta, k} = \{v_{\delta, k} \in C(\Omega_k) : v_{\delta, k}|_T \in P_1(T), \quad T \in \mathcal{T}_{h_k}, \quad v_{\delta, k}|_{\partial\Omega \cap \partial\Omega_k} = 0\},$$

and the global finite element spaces

$$X_\delta = \prod_{k=1}^K X_{\delta, k}, \quad \text{where} \quad \delta = (h_1, h_2, \dots, h_K).$$

Notice that the trace of the triangulation \mathcal{T}_{h_k} over Γ_{kl} , $1 \leq k \leq K$, $l \in \bar{k}$, with vertices $v_{1,kl}$ and $v_{2,kl}$ results in a regular triangulation denoted by t_{kl} , where \bar{k} is the class of the indices $l \in \bar{k}$ with $l > k$ and \bar{k} is denotes the set of all indices l so that kl exist. The trace space $W_{\delta, kl}$ of the functions in $X_{\delta, k}$ is given by (see Figure 2):

$$W_{\delta, kl} = \{\phi_{\delta, kl} \in C(\Gamma_{kl}) : t \in t_{kl}, \quad \phi_{\delta, kl} \in P_1(t)\},$$

the approximation of the local Lagrange multiplier is defined as

$$M_{\delta, kl} = \{\phi_{\delta, kl} \in W_{\delta, kl} : t \in t_{kl}, \quad \text{or} \quad \phi_{\delta, kl} \in P_0(t) \quad \text{if} \quad v_{1,kl} \quad \text{or} \quad v_{2,kl} \in t\},$$

and the global finite element space on the interface is

$$M_\delta = \prod_{k=1}^K \prod_{l \in \bar{k}} M_{\delta, kl}.$$

We now define a bilinear form on $X_\delta \times M_\delta$ by

$$B(v_\delta, \phi_\delta) = \sum_{k=1}^K \langle v_{\delta, k}, \phi_{\delta, k} \rangle_{*, \partial\Omega_k} = \sum_{k=1}^K \sum_{l \in \bar{k}} \int_{\Gamma_{kl}} \phi_{\delta, kl} (v_{\delta, k} - v_{\delta, l}) ds.$$

Thus, the mortar finite element approximation of the solution of (2) is defined by (see, e.g. [Bel97, BDM90, BMP94]):

$$\begin{aligned} A(u_\delta, v_\delta) + B(v_\delta, \psi_\delta) &= (f, v_\delta), \quad v_\delta \in X_\delta, \\ B(u_\delta, \phi_\delta) &= 0, \quad \phi_\delta \in M_\delta. \end{aligned} \quad (3)$$

If the space V_δ of nonconforming approximations of functions in $H_0^1(\Omega)$ is introduced by:

$$V_\delta = \{v_\delta \in X_\delta : B(v_\delta, \phi_\delta) = 0, \quad \phi_\delta \in M_\delta\},$$

then the problem (3) is equivalent to the problem of finding $u_\delta \in V_\delta$ such that

$$A(u_\delta, v_\delta) = (f, v_\delta) \quad v_\delta \in V_\delta. \quad (4)$$

Now we shall introduce the mortar finite volume element approximation of the model problem (3). For a given triangulation \mathcal{T}_{h_k} , we construct a dual mesh $\mathcal{T}_{h_k}^*$ based upon \mathcal{T}_{h_k} whose elements are called control volumes.

There are various ways of introducing regular control volume grids \mathcal{T}_{δ}^* . In the most popular control volume partitions, the mediacenter of the finite element T is connected with the midpoints of the edges of T . These types of volumes can be introduced for any finite element partition \mathcal{T}_{h_k} and leads to relatively simple calculations. If the vertex is on the interface Γ_{kl} , then ‘‘half’’ control volume (shaded regions in Figure 1) is used.

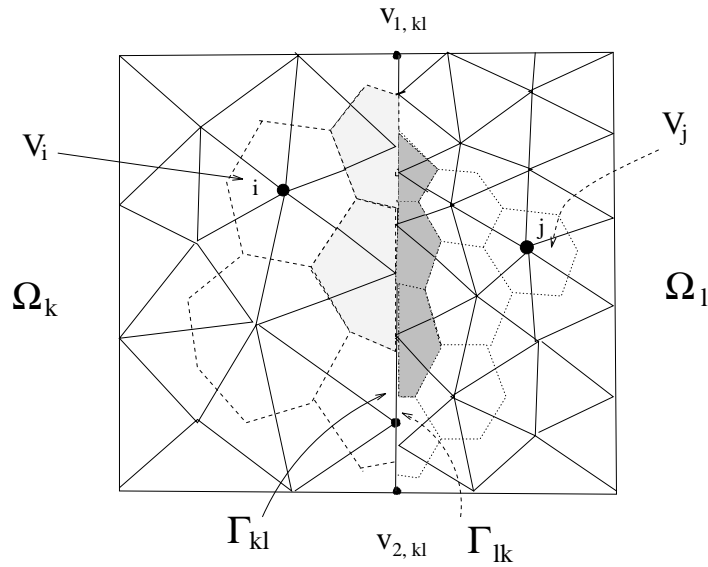


Figure 1 Interfaces Γ_{kl} and Γ_{lk} with $v_{1,kl}$ and $v_{2,kl}$ as two end points, triangulation \mathcal{T}_{h_k} and \mathcal{T}_{h_l} , and the volumes in Ω_k and Ω_l . The triangulation t_{kl} and t_{lk} are different on the interface due to non-matching grids.

For the finite element space X_{δ} we can define its dual volume element space $X_{\delta}^* = \prod_{k=1}^K X_{\delta,k}^*$, where

$$X_{\delta,k}^* = \{v_k \in L^2(\Omega_k) : v_k|_V \text{ is constant over } V \in \mathcal{T}_{h_k}^* \text{ and } v_k|_{\partial\Omega \setminus \partial\Omega_k} = 0\}.$$

Obviously, $X_{\delta,k}^* = \text{span}\{\chi_{i,k}(V) : V \in \mathcal{T}_{h_k}^*\}$, where $\chi_{i,k}$ is the characteristic function of the volume $V_{i,k}$. Let $I_{h_k} : C(\Omega_k) \rightarrow X_{\delta,k}$ be the interpolation operator and $I_{h_k}^* : C(\Omega_k) \rightarrow X_{\delta,k}^*$ be the piecewise constant interpolation operator, that is

$$I_{h_k}^* u = \sum_{x_{i,k} \in N_{h_k}} u_{i,k} \chi_{i,k}(x), \text{ where } u_{i,k} = u(x_{i,k}).$$

Then we set $I_{\delta} = \prod_{k=1}^K I_{h_k}$ and $I_{\delta}^* = \prod_{k=1}^K I_{h_k}^*$. With the above preparation, we can combine the finite volume approximation (see, e.g. [Cai91, ELL98, LC94, Mis98])

with the mortar approach to define our *mortar finite volume element method*: find $(u_\delta, \psi_\delta) \in X_\delta \times M_\delta$ such that

$$\begin{aligned} A(u_\delta, I_\delta^* v_\delta) + B(v_\delta, \psi_\delta) &= (f, I_\delta^* v_\delta), & v_\delta \in X_\delta, \\ B(u_\delta, \phi_\delta) &= 0, & \phi_\delta \in M_\delta, \end{aligned} \quad (5)$$

where

$$\begin{aligned} A(u_\delta, I_\delta^* v_\delta) &= - \sum_{k=1}^K \sum_{j \in N_{h_k}} v_{j,k} \int_{\partial V_{j,k}} A(x) \nabla u_{\delta,k} \cdot \mathbf{n}_k ds, \\ (f, I_\delta^* v_\delta) &= \sum_{k=1}^K \sum_{j \in N_{h_k}} v_{j,k} \int_{V_{j,k}} f(x) dx. \end{aligned}$$

This problem is equivalent to the following problem: find $u_\delta \in V_\delta$ such that

$$A(u_\delta, I_\delta^* v_\delta) = (f, I_\delta^* v_\delta), \quad v_\delta \in V_\delta. \quad (6)$$

Remark: We keep the same piecewise linear element spaces on the interfaces and formulate our mortar finite volume approximations only on the subdomains. This alone in fact is enough to preserve the basic feature of finite volume element method, that is, both (5) and (6) are locally conservative. The weak compatibility condition of the spaces X_δ and M_δ are satisfied automatically:

$$\{\phi_\delta : B(v_\delta, \phi_\delta) = 0, \quad \forall v_\delta \in X_\delta\} = \{0\}, \quad (7)$$

which guarantees that there is no spurious modes generated for the normal derivatives of the solution using this discretization. In other words, our mortar finite volume element formulation has the nice properties of mortar finite element method.

In [ELLL98] we have introduced another formulation of the mortar finite volume element method with piecewise constant volume element approximation on the interfaces. The stability, convergence and error estimates for this type of scheme can be obtained in the framework presented above. Similarly, methods for geometrically nonconforming subdomains or overlapping domains as those shown on Figure 24 can be introduced as well (see [ELLL98]).

Error Estimates

The following error estimate has been proved in [ELLL98]:

Theorem 1 *Assume that \mathcal{T}_δ is regular, then the unique solution pair $(u_\delta, \psi_\delta) \in X_\delta \times M_\delta$ exists for the finite volume element mortar formulation and satisfies the error estimates:*

$$\|u - u_\delta\|_X \leq C \sum_{k=1}^K h_k \|u\|_H^2(\Omega_k) + C \sum_{k=1}^K h_k \|f\|_{L^2(\Omega)}.$$

A similar estimate is valid for the M -norm of the error in the Lagrange multipliers as well.

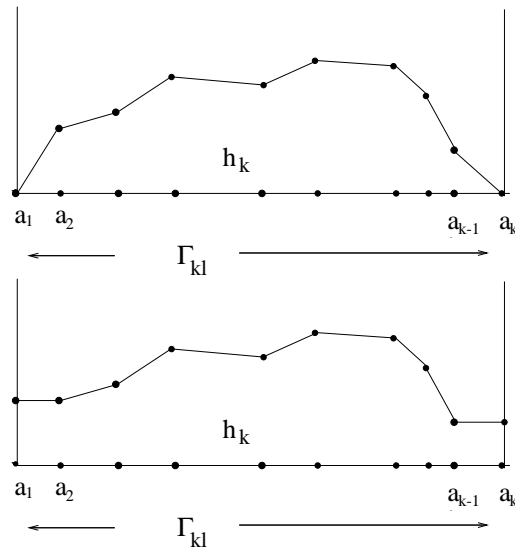


Figure 2 Left: a function from $W_{\delta,kl}$; right: a function from $M_{\delta,kl}$.

REFERENCES

- [AY97] Arbogast T. and Yotov I. (1997) A non-mortar mixed finite element method for elliptic problems on non-matching multi-block grids. *Computer Methods in Applied Mechanics and Engineering* 149: 255–265.
- [BDM90] Bernardi C., Debit N., and Maday Y. (1990) Coupling finite element and spectral methods: First results. *Math. Comput.* 54: 21–39.
- [Bel97] Belgacem F. (1997) The mortar finite element method with Lagrange multiplier. preprint.
- [BM94] Belgacem F. and Maday Y. (1994) Nonconforming spectral element methodology tuned to parallel implementation. *Computer Methods in Applied Mechanics and Engineering* 116: 59–67.
- [BM97] Bernardi C. and Maday Y. (1997) The mortar element method for three dimensional finite elements. *M2NA* 31: 269–301.
- [BMP94] Bernardi C., Maday Y., and Patera A. (1994) *A new nonconforming approach to domain decomposition: the mortar element method.* 299.
- [BMSL89] Bernardi C., Maday Y., and Sacchi-Landriani G. (1989) Nonconforming matching conditions for coupling spectral finite element methods. *Applied Numer. Math.* 54: 64–84.
- [BV96] Benkhaldoun F. and Vilsmeier R. (eds) (1996) *Proc. Int. Conference on Finite Volumes for Complex Applications.*
- [Cai91] Cai Z. (1991) On the finite volume element method. *Numer. Math.* 58: 713–735.
- [Cas] Casarin M. Schwarz preconditioners for spectral and mortar finite element methods with applications to incompressible fluids. Technical report, Department of Computer Sciences.
- [CDS99] Cai X.-C., Dryja M., and Sarkis M. (1999) Overlapping nonmatching grid mortar element methods for elliptic problems. *SIAM J. Numer. Anal.* 35:

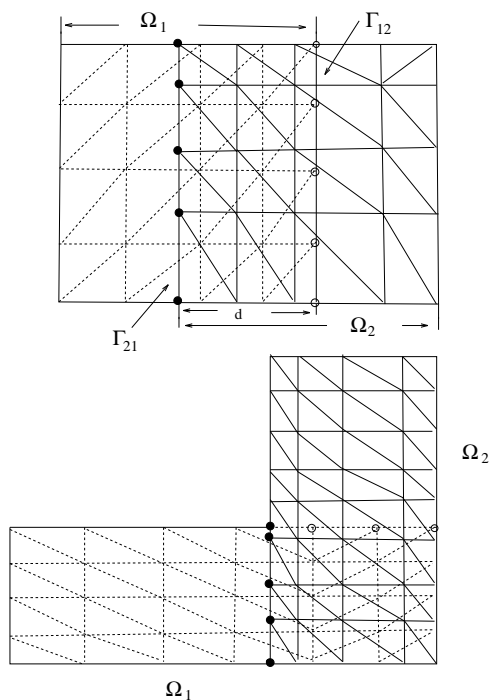


Figure 3 Overlapping subdomains.

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- [ELL98] Ewing R., Lazarov R., and Lin Y. (1998) Finite volume element approximations of nonlocal reactive flows in porous media. Technical report, Texas A&M University, Institute for Scientific Computation.
- [ELLL98] Ewing R., Lazarov R., Lin T., and Lin Y. (1998) The mortar finite volume element methods and domain decomposition. Technical report, Texas A&M University, Institute for Scientific Computation.
- [LC94] Li R. and Chen Z. (1994) *The generalized difference method for differential equations*. Jilin University Publishing House.
- [Mis98] Mishev I. (1998) Finite volume methods on Voronoi meshes. *Numer. Methods for Partial Differential Equations* 14: 193–212.