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Talk: Time- and Space-Decomposition Methods for Parabolic Problems and Applications in Multiphysics Problems.

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Outline of the talk

- 1.) Introduction
- 2.) Decomposition-methods
- 3.) Time- and Space Decomposition methods
- 4.) Combined Time- and Space Decomposition Methods
- 5.) A priori Error-Estimate
- 6.) Numerical experiments
- 7.) Future Works

Motivation and Ideas

Decoupling of multi physics problems to simpler physics problems

Embedding the physical characteristics to the numerical methods (conservation of physics)

Parallelization and accelerating the solver-process

Higher order methods for time and space

Methods for non-smooth and degenerate problems

Fast computations for complicate and decoupable problems

Model-Equation

Systems of parabolic-differential equations with first order timederivation and second order spatial-derivations

$$\begin{aligned} \frac{\partial c}{\partial t} &= f(c) + Ac + Bc , \text{ in } \Omega \times (0, T) , \\ c(x, t) &= g(x, t) , \text{ on } \partial \Omega \times (0, T) \text{ (Boundary-Condition)} , \\ c(x, 0) &= c_0(x) , \text{ in } \Omega \text{ (Initial-Condition)} , \end{aligned}$$

where
$$c = (c_1, ..., c_n)^t$$
 and $f(c) = (f_1(c), ..., f_n(c))^t$,

$$A = \begin{pmatrix} -v_{11} \cdot \nabla & \cdots & -v_{n1} \cdot \nabla \\ \cdots & \cdots & \cdots \\ -v_{1n} \cdot \nabla & \cdots & -v_{nn} \cdot \nabla \end{pmatrix}, B = \begin{pmatrix} \nabla D_{11} \cdot \nabla & \cdots & \nabla D_{n1} \cdot \nabla \\ \cdots & \cdots & \cdots \\ \nabla D_{1n} \cdot \nabla & \cdots & \nabla D_{nn} \cdot \nabla \end{pmatrix},$$

Convection- and diffusion-operator with $A, B : X \to X$ and $X = \mathbb{R}^n$ a matrix-space. sufficient smoothness $c_i \in C^{2,1}(\Omega, [0, T])$ for i = 1, ..., n

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Decomposition Methods

Ideas :

Decoupling the time-scales, space-scales. Decoupling the multi-physics. Time-adaptivity, Space-adaptivity. Parallelization in Time and Space.

Methods :

Operator-Splitting and Variational Splitting Methods (Time). Iterative and extended Operator Splitting Methods (Time). Waveform-Relaxation-Methods (Time). Schwarz Wave form relaxation method (Space). Additive and Multiplicative Schwarz method (Space). Partition of Units combined with Splitting methods (Time and Space).

Spatial decomposition method : Overlapping Schwarz wave form relaxation method

Given the following model problem

$$u_t + Lu = f , \text{ in } \Omega \times (0, T) , \qquad (2)$$

$$\overline{\Omega} \times (0, T) := \overline{\Omega}_1 \times (0, T) \cup \overline{\Omega}_2 \times (0, T) ,$$

$$u(x, 0) = u_0 , \text{ (Initial-Condition)} ,$$

$$u = g , \text{ on } \partial\Omega \times (0, T) ,$$

where L denotes for each time t a second-order partial differential operator $Lu = -\nabla D \nabla u + v \nabla u + cu$ for the given coefficients $D \in \mathbb{R}^+, v \in \mathbb{R}^n, c \in \mathbb{R}^+$, and n is the dimension of the space.

Schwarz-Waveform Relaxation method

We consider the method for two half steps, associated with the two subdomains and we solve 2 subproblems

$$u_{1t} + Lu_1^n = f , \text{ in } \Omega_1 \times (0, T) , \qquad (3)$$

$$u_1(x, 0) = u_{10} , \text{ (Initial-Condition)} ,$$

$$u_1^n = g , \text{ on } L_0 = \partial\Omega \times (0, T) \cap \partial\Omega_1 \times (0, T) ,$$

$$u_1^n = u_2^{n-1} , \text{ on } L_2 = \partial\Omega_1 \times (0, T) \setminus \partial\Omega \times (0, T) ,$$

$$u_{2t} + Lu_2^n = f , \text{ in } \Omega_2 \times (0, T) , \qquad (4)$$

$$u_2(x, 0) = u_{20} , \text{ (Initial-Condition)} ,$$

$$u_2^n = g , \text{ on } L_3 = \partial\Omega \times (0, T) \cap \partial\Omega_2 \times (0, T) ,$$

$$u_2^n = u_1^n , \text{ on } L_1 = \partial\Omega_2 \times (0, T) \setminus \partial\Omega \times (0, T); ,$$

Error of an Overlapping Schwarz wave form relaxation for the scalar convection reaction diffusion equation

We consider the convection diffusion reaction equation, given by

$$u_t = D u_{xx} - \nu u_x - \lambda u , \qquad (5)$$

defined on the domain $\Omega = [0, L]$ for $T = [T_0, T_f]$, with the following initial and boundary conditions

$$u(0,t) = f_1(t), \quad u(L,t) = f_2(t), \quad u(x,T_0) = u_0$$

To solve the model problem using overlapping Schwarz wave form relaxation method, we subdivide the domain Ω in two overlapping sub-domains $\Omega_1 = [0, L_2]$ and $\Omega_2 = [L_1, L]$, where $L_1 < L_2$ and $\Omega_1 \bigcap \Omega_2 = [L_1, L_2]$ is the overlapping region for Ω_1 and Ω_2 .

The convergence and error-estimates of $e^{k+1} = u - u_1^{k+1}$ and $d^{k+1} = u - u_2^{k+1}$ given by (3) and (4) respectively, are presented by

the following theorem

Theorem 1. Let e^{k+1} and d^{k+1} be the error from the solution of the subproblems (3) and (4) by Schwarz wave form relaxation over Ω_1 and Ω_2 , respectively, then

$$||e^{k+2}(L_1,t)||_{\infty} \le \gamma ||e^k(L_1,t)||_{\infty}$$

and

$$||d^{k+2}(L_2,t)||_{\infty} \le \gamma ||d^k(L_1,t)||_{\infty}$$
,

where

$$\gamma = \frac{\sinh(\beta L_1)}{\sinh(\beta L_2)} \frac{\sinh(\beta(L_2 - L))}{\sinh(\beta(L_1 - L))} < 1 ,$$
with $\beta = \frac{\sqrt{\nu^2 + 4D\lambda}}{2D}$.

Proof see [Geiser & Daoud, in review to NMPDE, 2006]

Time-Decomposition methods : Sequential Splitting methods

Idea: Decoupling of complex equations in simpler equations, solving simpler equations and re-coupling the results over the initial-conditions.

Equations: $\partial_t c = Ac + Bc$, where the initial-conditions are $c(t^n) = c^n$, (or Variational-formulation: $(\partial_t c, v) = (Ac, v) + (Bc, v)$.)

Splitting-method of first order

$$\partial_t c^* = Ac^*$$
 with $c^*(t^n) = c^n$,
 $\partial_t c^{**} = Bc^{**}$ with $c^{**}(t^n) = c^*(t^{n+1})$,

where the results of the methods are $c(t^{n+1}) = c^{**}(t^{n+1})$, and there are some splitting-errors for these methods, Literature : [Strang 68], [Karlsen et al 2001].

Splitting-Errors of the Method

The error of the splitting-method of first order is

$$\partial_t c = (B+A)c$$
,
 $\tilde{c} = \exp(\tau(B+A))c(t^n)$

Local error for the decomposition and the full solution

$$e(c) = \tilde{c}(t^{n} + \tau) - \exp(\tau B) \exp(\tau A)c(t^{n}) ,$$

= $\exp(\tau (B + A))c(t^{n}) - \exp(\tau B) \exp(\tau A)c(t^{n}) ,$
 $e(c)/\tau = \frac{1}{2}\tau (BA - AB)c(t^{n}) + O(\tau^{2}) ,$

 $O(\tau)$ for A, B not commuting, otherwise one get exact results, where $\tau = t^{n+1} - t^n$, [Strang 68].

Higher order splitting-methods

Strang or Strang-Marchuk-Splitting, cf. [Marchuk 68, Strang68]

$$\begin{aligned} \frac{\partial c^*(t)}{\partial t} &= Ac^*(t), \text{ with } t^n \leq t \leq t^{n+1/2} \text{ and } c^*(t^n) = c_{\rm sp}^n, \end{aligned} \tag{6} \\ \frac{\partial c^{**}(t)}{\partial t} &= Bc^{**}(t), \text{ with } t^n \leq t \leq t^{n+1}, c^{**}(t^n) = c^*(t^{n+1/2}), \\ \frac{\partial c^{***}(t)}{\partial t} &= Ac^{***}(t), t^{n+1/2} \leq t \leq t^{n+1}, c^{***}(t^{n+1/2}) = c^{**}(t^{n+1}), \end{aligned}$$

where $t^{n+1/2} = t^n + 0.5\tau_n$ and the approximation on the next time level t^{n+1} is defined as $c_{sp}^{n+1} = c^{***}(t^{n+1})$.

The splitting error of the Strang splitting is

$$\rho_n = \frac{1}{24} \tau_n^2([B, [B, A]] - 2[A, [A, B]]) \ c(t^n) + O(\tau_n^3) \ , \quad (7)$$

see, e.g.[Hundsdorfer, Verwer 2003].

Combined Methods

Introduction Iterative splitting-Methods

$$\frac{\partial c_i(t)}{\partial t} = Ac_i(t) + Bc_{i-1}(t), \text{ with } c_i(t^n) = c_{sp}^n, \quad (8)$$
$$\frac{\partial c_{i+1}(t)}{\partial t} = Ac_i(t) + Bc_{i+1}(t), \text{ with } c_{i+1}(t^n) = c_{sp}^n, \quad (9)$$

where $c_0(t)$ is any fixed function for each iteration. (Here, as before, $c_{\rm sp}^n$ denotes the known split approximation at the time level $t = t^n$.) The split approximation at the time-level $t = t^{n+1}$ is defined as $c_{\rm sp}^{n+1} = c_{2m+1}(t^{n+1})$. (Clearly, the functions $c_k(t)$ (k = i - 1, i, i + 1) depend on the interval $[t^n, t^{n+1}]$, too, but, for the sake of simplicity, in our notation we omit the dependence on n.)

Error for the Iterative splitting-method

Theorem 2. The error for the splitting methods is given as :

$$||e_i|| = K||B||\tau_n||e_{i-1}|| + O(\tau_n^2)$$
(10)
and hence

$$||e_{2m+1}|| = K_m ||e_0||\tau_n^{2m} + O(\tau_n^{2m+1}),$$
(11)

where τ_n is the time-step, e_0 the initial error $e_0(t) = c(t) - c_0(t)$ and m the number of iteration-steps, K and K_m are constants, ||B|| is the maximum norm of operator B and A and B are bounded, monotone operators.

Proof : Taylor-expansion and estimation of \exp -functions. See the work Geiser, Farago (2005).

The combined time-space iterative splitting method

Based on the iterative operator-splitting method we extend the splitting method be an embedded Schwarz-waveform-relaxation method.

We solve the following sub-problems consecutively for $i = 0, 2, \ldots 2m$ and $j = 0, 2, \ldots 2n$. In this notation i represents the iteration index for the time-splitting and j represents the iteration index for the spatial-splitting.

Initial idea:

$$\begin{aligned} \frac{\partial c_{i,j}(t)}{\partial t} &= A|_{\Omega_1} c_{i,j}(t) + A|_{\Omega_2} c_{i,j-1}(t) + B|_{\Omega_1} c_{i-1,j}(t) + B|_{\Omega_2} c_{i-1,j-1}(t), \end{aligned}$$
(12)

$$\begin{aligned} \frac{\partial c_{i+1,j}(t)}{\partial t} &= A|_{\Omega_1} c_{i,j}(t) + A|_{\Omega_2} c_{i,j-1}(t) + B|_{\Omega_1} c_{i+1,j}(t) + B|_{\Omega_2} c_{i-1,j-1}(t), \end{aligned}$$
(13)

$$\begin{aligned} \frac{\partial c_{i,j+1}(t)}{\partial t} &= A|_{\Omega_1} c_{i,j}(t) + A|_{\Omega_2} c_{i,j+1}(t) + B|_{\Omega_1} c_{i+1,j}(t) + B|_{\Omega_2} c_{i-1,j-1}(t), \end{aligned}$$
(14)

$$\begin{aligned} \frac{\partial c_{i+1,j+1}(t)}{\partial t} &= A|_{\Omega_1} c_{i,j}(t) + A|_{\Omega_2} c_{i,j+1}(t) + B|_{\Omega_1} c_{i+1,j}(t) + B|_{\Omega_2} c_{i+1,j+1}(t), \end{aligned}$$
(15)

where c^n is the known split approximation at the time level $t = t^n$.

The nonoverlapping time-space iterative splitting method

We denote for the semi-discretisation in space the variable k as the node for the point x_k and we obtain $k \in (0, \ldots, p)$, where p is the number of nodes. We have the decomposition if the space, where Ω_1 is of the points $0, \ldots, p/2$ and Ω_2 is of $p/2 + 1, \ldots, p$, we assume p is even. So we assume $\Omega_1 \cap \Omega_2 = \{\}$ and we have the following algorithm :

$$\frac{\partial(c_{i,j})_{k}(t)}{\partial t} = \tilde{A}|_{\Omega_{1}}(c_{i,j})_{k}(t) + \tilde{A}|_{\Omega_{2}}(c_{i,j-1})_{k}(t) + \tilde{B}|_{\Omega_{1}}(c_{i-1,j})_{k}(t) + \tilde{B}|_{\Omega_{2}}(c_{i-1,j-1})_{k}(t), \quad (16)$$
with $(c_{i,j})_{k}(t)(t^{n}) = (c^{n})_{k}$

$$\frac{\partial(c_{i+1,j})_{k}(t)}{\partial t} = \tilde{A}|_{\Omega_{1}}(c_{i,j})_{k}(t) + \tilde{A}|_{\Omega_{2}}(c_{i,j-1})_{k}(t) + \tilde{B}|_{\Omega_{1}}(c_{i+1,j})_{k} + \tilde{B}|_{\Omega_{2}}(c_{i-1,j-1})_{k}(t), \quad (17)$$
with $(c_{i+1,j})_{k}(t^{n}) = (c^{n})_{k}$

$$\frac{\partial(c_{i,j+1})_{k}(t^{n})(t)}{\partial t} = \tilde{A}|_{\Omega_{1}}(c_{i,j})_{k}(t^{n})(t) + \tilde{A}|_{\Omega_{2}}(c_{i,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{1}}(c_{i+1,j})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{2}}(c_{i-1,j-1})_{k}(t^{n})(t), \quad (18)$$

$$\frac{\partial(c_{i+1,j+1})_{k}(t^{n})(t)}{\partial t} = \tilde{A}|_{\Omega_{1}}(c_{i,j})_{k}(t^{n})(t) + \tilde{A}|_{\Omega_{2}}(c_{i,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{1}}(c_{i+1,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{2}}(c_{i+1,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{2}}(c_{i+1,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{2}}(c_{i+1,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{2}}(c_{i+1,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{2}}(c_{i+1,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{2}}(c_{i+1,j+1})_{k}(t^{n})(t) + \tilde{B}|_{\Omega_{2}}(c_{i+1,j+1})_{k}(t^{n})(t), \quad (19)$$

where c^n is the known split approximation at the time level $t = t^n$.

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We have the operators :

$$\tilde{A}|_{\Omega_1}(c_{i,j})_k = \begin{cases} A(c_{i,j})_k & \text{for } k \in \{0, \dots, p/2\} \\ 0 & \text{for } k \in \{p/2+1, \dots, p\} \end{cases}$$
(20)

$$\tilde{A}|_{\Omega_2}(c_{i,j})_k = \begin{cases} 0 & \text{for } k \in \{0, \dots, p/2\} \\ A(c_{i,j})_k & \text{for } k \in \{p/2, \dots, p\} \end{cases}$$
(21)

Similar are the assignments for operator B.

$$\tilde{B}|_{\Omega_1}(c_{i,j})_k = \begin{cases} B(c_{i,j})_k & \text{for } k \in \{0, \dots, p/2\} \\ 0 & \text{for } k \in \{p/2+1, \dots, p\} \end{cases}$$
(22)

$$\tilde{B}|_{\Omega_2}(c_{i,j})_k = \begin{cases} 0 & \text{for } k \in \{0, \dots, p/2\} \\ B(c_{i,j})_k & \text{for } k \in \{p/2, \dots, p\} \end{cases}$$
(23)

The overlapping time-space iterative splitting method

We denote for the semi-discretisation in space the variable k as the node for the point x_k and we obtain $k \in (0, \ldots, p)$, where p is the number of nodes. Now we assume the overlapping case, so we assume $\Omega_1 \cap \Omega_2 \neq \{\}$. We have the following sets : $\Omega \setminus \Omega_2 = \{0, \ldots, p_1\}$, $\Omega_1 \cap \Omega_2 = \{p_1 + 1, \ldots, p_2\}$ and $\Omega \setminus \Omega_1 = \{p_2 + 1, \ldots, p\}$. We assume $p_1 < p_2 < p$ and can derive the following overlapping algorithm :

$$\begin{aligned} \frac{\partial(c_{i,j})_{k}(t)}{\partial t} &= \tilde{A}|_{\Omega \setminus \Omega_{2}}(c_{i,j})_{k}(t) + \tilde{A}|_{\Omega_{1} \cap \Omega_{2}}(c_{i,j}, c_{i,j-1})_{k}(t) + \tilde{A}|_{\Omega \setminus \Omega_{1}}(c_{i,j-1})_{k}(t) \\ &+ \tilde{B}|_{\Omega \setminus \Omega_{2}}(c_{i-1,j})_{k}(t) + \tilde{B}|_{\Omega_{1} \cap \Omega_{2}}(c_{i-1,j}, c_{i-1,j-1})_{k}(t) + \tilde{B}|_{\Omega \setminus \Omega_{1}}(c_{i-1,j-1})_{k}(t), \\ &\text{with } (c_{i,j})_{k}(t)(t^{n}) &= (c^{n})_{k} \end{aligned} \tag{24} \\ &\frac{\partial(c_{i+1,j})_{k}(t)}{\partial t} &= \tilde{A}|_{\Omega \setminus \Omega_{2}}(c_{i,j})_{k}(t) + \tilde{A}|_{\Omega_{1} \cap \Omega_{2}}(c_{i,j}, c_{i,j-1})_{k}(t) + \tilde{A}|_{\Omega \setminus \Omega_{1}}(c_{i,j-1})_{k}(t) \\ &+ \tilde{B}|_{\Omega \setminus \Omega_{2}}(c_{i+1,j})_{k}(t) + \tilde{B}|_{\Omega_{1} \cap \Omega_{2}}(c_{i+1,j}, c_{i-1,j-1})_{k}(t) + \tilde{B}|_{\Omega \setminus \Omega_{1}}(c_{i-1,j-1})_{k}(t), \\ &\text{with } (c_{i+1,j})_{k}(t^{n}) &= (c^{n})_{k} \end{aligned} \tag{25} \\ &\frac{\partial(c_{i,j+1})_{k}(t)}{\partial t} &= \tilde{A}|_{\Omega \setminus \Omega_{2}}(c_{i,j})_{k}(t) + \tilde{A}|_{\Omega_{1} \cap \Omega_{2}}(c_{i,j+1}, c_{i,j})_{k}(t) + \tilde{A}|_{\Omega \setminus \Omega_{1}}(c_{i,j+1})_{k}(t) \\ &+ \tilde{B}|_{\Omega \setminus \Omega_{2}}(c_{i+1,j})_{k}(t) + \tilde{B}|_{\Omega_{1} \cap \Omega_{2}}(c_{i+1,j}, c_{i-1,j-1})_{k}(t) + \tilde{B}|_{\Omega \setminus \Omega_{1}}(c_{i-1,j-1})_{k}(t), \\ &\text{with } (c_{i,j+1})_{k}(t^{n})(t^{n}) &= (c^{n})_{k}(t^{n}) \end{aligned} \tag{26} \\ &\frac{\partial(c_{i+1,j+1})_{k}(t)}{\partial t} &= \tilde{A}|_{\Omega \setminus \Omega_{2}}(c_{i,j})_{k}(t) + \tilde{A}|_{\Omega_{1} \cap \Omega_{2}}(c_{i,j+1}, c_{i,j})_{k}(t) + \tilde{A}|_{\Omega \setminus \Omega_{1}}(c_{i,j+1})_{k}(t) \\ &+ \tilde{B}|_{\Omega \setminus \Omega_{2}}(c_{i+1,j})_{k}(t) + \tilde{B}|_{\Omega_{1} \cap \Omega_{2}}(c_{i+1,j}, c_{i+1,j+1})_{k}(t) + \tilde{A}|_{\Omega \setminus \Omega_{1}}(c_{i,j+1})_{k}(t) \end{aligned} \end{aligned}$$

We have the operators :

$$\tilde{A}|_{\Omega \setminus \Omega_2}(c_{i,j})_k = \begin{cases} A(c_{i,j})_k & \text{for } k \in \{0, \dots, p_1\} \\ 0 & \text{for } k \in \{p_1 + 1, \dots, p\} \end{cases}$$
(28)

$$\tilde{A}|_{\Omega_1 \cap \Omega_2}(c_{i,j}, c_{i,j+1})_k = \begin{cases} A((c_{i,j} + c_{i,j+1})/2)_k & \text{for } k \in \{p_1 + 1, \dots, p_2\} \\ 0 & \text{for } k \in \{p_2 + 1, \dots, p\} \end{cases}$$
(29)

$$\tilde{A}|_{\Omega \setminus \Omega_1}(c_{i,j})_k = \begin{cases} 0 & \text{for } k \in \{0, \dots, p_2\} \\ A(c_{i,j})_k & \text{for } k \in \{p_2 + 1, \dots, p\} \end{cases}$$
(30)

Similar are the assignments for operator B.

$$\tilde{B}|_{\Omega\setminus\Omega_2}(c_{i,j})_k = \begin{cases} B(c_{i,j})_k & \text{for } k \in \{0,\dots,p_1\}\\ 0 & \text{for } k \in \{p_1+1,\dots,p\} \end{cases}$$
(31)

$$\tilde{B}|_{\Omega_1 \cap \Omega_2}(c_{i,j}, c_{i,j+1})_k = \begin{cases} B((c_{i,j} + c_{i,j+1})/2)_k & \text{for } k \in \{p_1 + 1, \dots, p_2\} \\ 0 & \text{for } k \in \{p_2 + 1, \dots, p\} \end{cases}$$
(32)

$$\tilde{B}|_{\Omega \setminus \Omega_1}(c_{i,j})_k = \begin{cases} 0 & \text{for } k \in \{0, \dots, p_2\} \\ B(c_{i,j})_k & \text{for } k \in \{p_2 + 1, \dots, p\} \end{cases}$$
(33)

Dicretisation of the operators

The discretization of the operators is given as :

$$A(c_{i,j})_k = D/(\Delta x)^2 (-(c_{i,j})_{k+1} + 2(c_{i,j})_k - (c_{i,j})_{k-1}) -v/\Delta x ((c_{i,j})_k - (c_{i,j})_{k-1})$$
(34)

$$B(c_{i,j})_k = \lambda(c_{i,j})_k , \qquad (35)$$

Consistency and stability analysis of the combined method

Theorem 3. Let us consider the nonlinear operator-equation in a Banach space **X**

$$\partial_t c(t) = A_1(c(t)) + A_2(c(t)) + B_1(c(t)) + B_2(c(t)), \quad 0 < t \le T ,$$

$$c(0) = c_0 , \qquad (36)$$

where $A_1, A_2, B_1, B_2, A_1 + A_2 + B_1 + B_2 : \mathbf{X} \to \mathbf{X}$ are given linear operators being generators of the C_0 -semigroup and $c_0 \in \mathbf{X}$ is a given element. Then the iteration process (12)–(15) is convergent and the rate of the convergence is of second order.

We obtain the iterative result :

we obtain

$$||e_{i,j}|| = K\tau_n ||e_{i-1,j-1}|| + \mathcal{O}(\tau_n^2),$$
(37)

and hence

$$\|e_{i+1,j+1}\| = K_1 \tau_n^2 \|e_{i-1,j-1}\| + \mathcal{O}(\tau_n^3),$$
(38)

which proves our statement.

Proof see [Geiser & Kravvaritis 2006]

Let us consider the iteration (12)–(15) on the sub-interval $[t^n, t^{n+1}]$. For the error function $e_i(t) = c(t) - c_i(t)$ we have the relations

$$\partial_{t} e_{i,j}(t) = A_{1}(e_{i,j}(t)) + A_{2}(e_{i,j-1}(t))$$

$$+ B_{1}(e_{i-1,j}(t)) + B_{2}(e_{i-1,j-1}(t)),$$

$$t \in (t^{n}, t^{n+1}], \ e_{i,j}(t^{n}) = 0,$$
(40)

$$\partial_{t} e_{i+1,j}(t) = A_{1}(e_{i,j}(t)) + A_{2}(e_{i,j-1}(t))$$

$$+ B_{1}(e_{i+1,j}(t)) + B_{2}(e_{i-1,j-1}(t)),$$

$$t \in (t^{n}, t^{n+1}], \ e_{i+1,j}(t^{n}) = 0 ,$$
(41)
(41)
(41)

 $\quad \text{and} \quad$

$$\partial_{t} e_{i,j+1}(t) = A_{1}(e_{i,j}(t)) + A_{2}(e_{i,j+1}(t))$$

$$+ B_{1}(e_{i+1,j}(t)) + B_{2}(e_{i-1,j-1}(t)),$$

$$t \in (t^{n}, t^{n+1}], \ e_{i,j+1}(t^{n}) = 0,$$
(43)

 and

$$\partial_t e_{i,j}(t) = A_1(e_{i,j}(t)) + A_2(e_{i,j+1}(t))$$

$$+B_1(e_{i+1,j}(t)) + B_2(e_{i+1,j+1}(t)),$$

$$t \in (t^n, t^{n+1}], \ e_{i,j}(t^n) = 0,$$
(45)

for $i, j = 0, 2, 4, \ldots$, with $e_{0,0}(0) = 0$ and $e_{-1,0} = e_{0,-1} = e_{-1,-1}(t) = c(t)$.

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In the following we derive the linear system of equations. We use the notations \mathbf{X}^2 for the product space $\mathbf{X} \times \mathbf{X}$ enabled with the norm $||(u,v)|| = \max\{||u||, ||v||\}$ $(u,v \in \mathbf{X})$. The elements $\mathcal{E}_i(t)$, $\mathcal{F}_i(t) \in \mathbf{X}^2$ and the linear operator $\mathcal{A} : \mathbf{X}^2 \to \mathbf{X}^2$ are defined as follows

$$\mathcal{E}_{i,j}(t) = \begin{bmatrix} e_{i,j}(t) \\ e_{i+1,j}(t) \\ e_{i,j+1}(t) \\ e_{i+1,j+1}(t) \end{bmatrix}; \quad \mathcal{A} = \begin{bmatrix} A_1 & 0 & 0 & 0 \\ A_1 & A_2 & 0 & 0 \\ A_1 & A_2 & B_1 & 0 \\ A_1 & A_2 & B_1 & B_2 \end{bmatrix}, \quad (47)$$

$$\mathcal{F}_{i,j}(t) = \begin{bmatrix} A_2(e_{i,j-1}(t)) + B_1(e_{i-1,j}(t)) + B_2(e_{i-1,j-1}) \\ B_1(e_{i-1,j}(t)) + B_2(e_{i-1,j-1}) \\ B_2(e_{i-1,j-1}) \\ 0 \end{bmatrix}.$$
(48)

Then, using the notations (48), the relations (40)-(46) can be written in the form

$$\partial_t \mathcal{E}_{i,j}(t) = \mathcal{A}\mathcal{E}_{i,j}(t) + \mathcal{F}_{i,j}(t), \quad t \in (t^n, t^{n+1}],$$

$$\mathcal{E}_{i,j}(t^n) = 0.$$
(49)

Due to our assumptions, \mathcal{A} is a generator of the one-parameter C_0 semigroup $(\mathcal{A}(t))_{t\geq 0}$.

Hence using the variations of constants formula, the solution of the abstract Cauchy problem (49) with homogeneous initial condition can be written as

$$\mathcal{E}_{i,j}(t) = \int_{t^n}^t \exp(\mathcal{A}(t-s))\mathcal{F}_{i,j}(s)ds, \quad t \in [t^n, t^{n+1}].$$
(50)

Hence, using the denotation

$$\|\mathcal{E}_{i,j}\|_{\infty} = \sup_{t \in [t^n, t^{n+1}]} \|\mathcal{E}_{i,j}(t)\|$$
, (51)

We could estimate the right hand side $\mathcal{F}_i(t)$ and $\exp(\mathcal{A}(t))$ We could then estimate the $\mathcal{F}_i(t)$ as $||\mathcal{F}_{i,j}(t)|| \leq C||e_{i-1,j-1}||$.

 $\quad \text{and} \quad$

$$\int_{t^n}^{t^{n+1}} \|\exp(\mathcal{A}(t-s))\| ds \le K_{\omega}(t), \quad t \in [t^n, t^{n+1}],$$
(53)

and hence

$$K_{\omega}(t) \leq \frac{K}{\omega} \left(\exp(\omega \tau_n) - 1 \right) = K \tau_n + \mathcal{O}(\tau_n^2) , \qquad (54)$$

We obtain the a priori error-estimates

$$||e_{i,j}|| = K\tau_n ||e_{i-1,j-1}|| + \mathcal{O}(\tau_n^2) .$$
(55)

(52)

Parallelization of the Time-Decomposition method : Windowing

The idea for parallelization in time are the windowing, that the processors has an amount of time-steps to compute and to share the end-result of the computation as an initial-condition for the next processor.



Numerical Experiments

We consider the one-dimensional convection-reaction-diffusion equation

$$\partial_t u + v \partial_x u - \partial_x D \partial_x u = -\lambda u , \text{ in } \Omega \times (T_0, T_f) ,$$
 (56)

$$u(x,0) = u_{ex}(x,0) , \text{ (Initial-Condition)}, \tag{57}$$

$$u(x,t) = u_{ex}(x,t)$$
, on $\partial \Omega \times (T_0, T_f)$, (58)

where $\Omega \times [T_0, T_f] = [0, 150] \times [100, 10^5].$

The exact solution is given as

$$u_{ex}(x,t) = \frac{u_0}{2\sqrt{D\pi t}} \exp(-\frac{(x-vt)^2}{4Dt}) \exp(-\lambda t) .$$
 (59)

The initial condition and the Dirichlet boundary conditions are defined using the exact solution (59) at starting time $T_0 = 100$ and with $u_0 = 1.0$. We have $\lambda = 10^{-5}$, v = 0.001 and D = 0.0001.

First example : A-B splitting combined with Schwarz wave form relaxation method

In order to solve the model problem using overlapping Schwarz wave form relaxation method, we divide the domain Ω in two overlapping sub-domains $\Omega_1 = [0, L_2]$ and $\Omega_2 = [L_1, L]$, where $L_1 < L_2$, and $\Omega_1 \bigcap \Omega_2 = [L_1, L_2]$ is the overlapping region for Ω_1 and Ω_2 .

For the sequential operator splitting method (A-B splitting). For this purpose we divide each of these two equations in terms of the operators $A = D \frac{\partial^2}{\partial x^2} - \nu \frac{\partial}{\partial x}$ and $B = -\lambda$.

For the discretization of equation (6) we apply the finite-difference method for the spatial discretization and the implicite Euler method for the time discretization. We provide a variety of results for several sizes of space- and time-partition, and also for various overlap sizes.

Precisely, we treat the cases $h=1,\ 0.5,\ 0.25$ as spatial step-size, $\Delta t=5,\ 10,\ 20$ as time step.

The considered subdomains are $\Omega_1 = [0, 80]$ and $\Omega_2 = [70, 150]$, $\Omega_1 = [0, 60]$ and $\Omega_2 = [30, 150]$ and $\Omega_1 = [0, 100]$ and $\Omega_2 = [30, 150]$, with overlap sizes 10, 30 and 70, respectively.

Both the approximated and the exact solution are evaluated at the end-time $t = 10^5$. The errors given in Table 2 are the maximum errors that occurred over the whole space domain, i.e. they are calculated using the ∞ -norm for vectors.

time-step	err	err	err	err	err	err
$\Delta t = 5$	2.85e - 3	2.24e - 3	1.28e - 3	2.66e - 4	2.21e - 4	2.20e - 4
$\Delta t = 10$	3.94e - 3	2.61e - 3	2.56e - 3	3.03e - 4	3.02e - 4	3.01e - 4
$\Delta t = 20$	5.03e - 3	2.81e - 3	2.73e - 3	8.51e - 4	5.22e - 4	5.14e - 4
overlap	10	30	70	10	30	70
space-step		h = 1			h = 0.5	

Table 1: Error for the scalar convection diffusion reaction-equation using the Schwarz waveform relaxation method for three different sizes of overlapping 10,30 and 70.

time-step	err	err	err	
$\Delta t = 5$	2.09e - 5	1.99e - 5	1.97e - 5	
$\Delta t = 10$	4.55e - 5	4.34e - 5	4.29e - 5	
$\Delta t = 20$	8.10e - 4	5.66e - 4	4.88e - 4	
overlap	10	30	70	
space-step	h = 0.25			

Table 2: Error for the scalar convection diffusion reaction-equation using the Schwarz waveform relaxation method for three different sizes of overlapping 10,30 and 70.

Second Example : Combined method : Time-Space iterative operator splitting method

For the solution of (56) with the combined time-space iterative splitting method we divide again the equation in terms of the operators

$$A = D\frac{\partial^2}{\partial x^2} - \nu \frac{\partial}{\partial x}$$

and

 $B = -\lambda.$

The index $k = 0, 1, \ldots p$ is associated with the subdomains, i.e. for $k = 0, \ldots, p/2$ we are working on Ω_1 and for $k = p/2 + 1, \ldots, p$ on Ω_2 . For the first set of values for k we have actually only the effect of the restrictions of the operators A and B on Ω_1 . Similarly, the second set of values for k indicates the action of the restrictions of both operators on Ω_2 .

The indices i and j are related to the time- and space-discretization, respectively. For every $k = 0, \ldots, p/2$ and for every interval of the space-discretization we solve the appropriate problems on Ω_1 , for every interval of the time-discretization. Similarly for $k = p/2 + 1, \ldots, p$ on Ω_2 .

By a closer examination of the scheme (24)–(27), taking into account the definitions (32)–(23), we observe that the problems to be solved in the innermost loop are of the form $\partial_t c = Ac + Bc$, $c(x, t^n) = c^n$, where c appears with appropriate indices i and j.

Both the approximated and the exact solution are evaluated at the end-time $t = 10^5$. The errors given in the following tables are the maximum errors that occurred over the whole space domain, i.e. they are calculated using the ∞ -norm for vectors.

The results are given in Table 4.

time-step	err	err	err	err	err	err
$\Delta t = 5$	4.38e - 2	1.47e - 2	3.49e - 3	2.59e - 4	2.13e - 4	1.54e - 4
$\Delta t = 10$	5.12e - 2	2.26e - 2	7.46e - 3	2.45e - 4	2.22e - 4	2.15e - 4
$\Delta t = 20$	6.14e - 2	4.39e - 2	1.20e - 2	7.43e - 4	5.21e - 4	4.53e - 4
overlap	10	30	70	10	30	70
space-step		h = 1			h = 0.5	

Table 3: Error for the scalar convection diffusion reaction-equation using the Schwarz waveform relaxation method for three different sizes of overlapping 10,30 and 70.

time-step	err	err	err	
$\Delta t = 5$	7.23e - 6	6.49e - 6	8.29e - 6	
$\Delta t = 10$	3.49e - 5	3.47e - 5	3.37e - 5	
$\Delta t = 20$	5.23e - 4	5.42e - 4	3.21e - 4	
overlap	10	30	70	
space-step	h = 0.25			

Table 4: Error for the scalar convection diffusion reaction-equation using the Schwarz waveform relaxation method for three different sizes of overlapping 10,30 and 70.

Future Work

1. Theory for the Stability of the time-space iterative splitting methods.

- 2. Commutative, non-commutative theory : How to decouple
- 3. Degenerated problems and non-smooth problems
- 4. Numerical examples