
Numerical Solution of Linear Elliptic Problems with Robin Boundary Conditions by a Least-Squares/Fictitious Domain Method

Roland Glowinski¹ and Qiaolin He²

¹ Department of Mathematics, University of Houston, Houston, TX 77204, USA; Institute of Advanced Study, The Hong Kong University of Science and Technology, Kowloon, Hong Kong, angelarim@aol.com

² Department of Mathematics, The Hong Kong University of Science and Technology, Kowloon, Hong Kong, hq1aa@ust.hk

1 Introduction

Motivated by the numerical simulation of particulate flow with slip boundary conditions at the interface fluid/particles, our goal, in this publication, is to discuss a fictitious domain method for the solution of linear elliptic boundary value problems with Robin boundary conditions. The method is of the virtual control type and relies on a least-squares formulation making the problem solvable by a conjugate gradient algorithm operating in a well chosen control space. Numerical results are presented; they suggest optimal orders of convergence for the finite element implementation of our fictitious domain method. A (brief) history of fictitious domain methods can be found in, e.g., [[3], Chap. 8].

2 Formulation of the Boundary Value Problem

Let Ω and ω be two bounded domains of \mathbf{R}^d , such that $d \geq 1$ and $\bar{\omega} \subset \Omega$ (see Fig. 1). We denote by Γ and γ the boundaries of Ω and ω , respectively. The Robin–Dirichlet

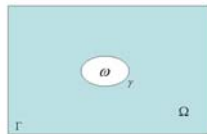


Fig. 1. Problem geometry.

problem under consideration reads as follows:

$$\begin{aligned} \alpha\psi - \mu\nabla^2\psi &= f \quad \text{in } \Omega \setminus \bar{\omega}, \\ \psi &= g_0 \quad \text{on } \Gamma, \quad \mu \left(\frac{\partial\psi}{\partial n} + \frac{\psi}{l_s} \right) = g_1 \quad \text{on } \gamma, \end{aligned} \tag{1}$$

where: $\alpha \geq 0, \mu > 0, f \in L^2(\Omega \setminus \bar{\omega}), g_0 \in H^{3/2}(\Gamma), g_1 \in H^{1/2}(\gamma), \mathbf{n}$ is the unit normal vector at γ pointing outward of $\Omega \setminus \bar{\omega}$ and l_s is a characteristic distance. We assume that Ω is convex and that γ is smooth. Problem (1) has a unique solution in $H^2(\Omega \setminus \bar{\omega})$ which is also the solution of the following linear variational problem:

$$\begin{aligned} \psi &\in H^1(\Omega \setminus \bar{\omega}), \psi = g_0 \quad \text{on } \Gamma, \\ \alpha \int_{\Omega \setminus \bar{\omega}} \psi \varphi dx + \mu \int_{\Omega \setminus \bar{\omega}} \nabla\psi \cdot \nabla\varphi dx + \frac{\mu}{l_s} \int_{\gamma} \psi \varphi d\gamma \\ &= \int_{\Omega \setminus \bar{\omega}} f \varphi dx + \int_{\gamma} g_1 \varphi d\gamma, \quad \forall \varphi \in V_0, \end{aligned} \tag{2}$$

where $dx = dx_1 \dots dx_d$ and $V_0 = \{\varphi | \varphi \in H^1(\Omega \setminus \bar{\omega}), \varphi = 0 \quad \text{on } \Gamma\}$.

3 A Least-Squares/Fictitious Domain Method for the Solution of Problem (1), (2)

3.1 A Fictitious Domain Formulation of Problem (1), (2)

We proceed as follows to define a fictitious domain variant of problem (1), (2):

(i) With $v \in L^2(\omega)$ we associate $\tilde{f}(v)$ defined by

$$\tilde{f}(v) \in L^2(\Omega), \quad \tilde{f}(v)|_{\Omega \setminus \bar{\omega}} = f, \quad \tilde{f}(v)|_{\omega} = v, \tag{3}$$

and then $\{\psi_1, \psi_2\}$ solution of the following elliptic system:

$$\alpha\psi_1 - \mu\nabla^2\psi_1 = \tilde{f}(v) \quad \text{in } \Omega, \quad \psi_1 = g_0 \quad \text{on } \Gamma, \tag{4}$$

$$\alpha\psi_2 - \mu\nabla^2\psi_2 = v \quad \text{in } \omega, \quad \mu \frac{\partial\psi_2}{\partial n} = \frac{\mu}{l_s} \psi_1 - g_1 \quad \text{on } \gamma. \tag{5}$$

Both problems (4) and (5) have a unique solution in $H^1(\Omega)$ and $H^1(\omega)$, respectively (actually, ψ_1 and ψ_2 have both the H^2 -regularity).

(ii) We define $\mathbf{A} : L^2(\omega) \rightarrow H^1(\omega)$ by

$$\mathbf{A}(v) = (\psi_2 - \psi_1)|_{\omega}. \tag{6}$$

Operator \mathbf{A} is clearly affine and continuous.

(iii) We observe that if v verifies $\mathbf{A}(v) = 0$, we then have $\psi_2 = \psi_1$ on ω and it is easy to see that the H^2 -regularity of ψ_1 and ψ_2 implies that $\psi_1|_{\Omega \setminus \bar{\omega}} = \psi$, where ψ is the solution of problem (1), (2). The problem is now to solve the functional equation

$$\mathbf{A}(u) = 0. \tag{7}$$

Indeed, the functional Eq. (7) has infinitely many solutions, but among these solutions only one is of minimal norm in $L^2(\omega)$.

Remark 1. Problem (7) can be viewed as an exact controllability problem in the sense of [2] (and as a virtual control problem in the sense of [4]). If a conjugate gradient algorithm is applied to a least-squares variant of (7) starting with 0 as initial guess, we have convergence to the unique solution of problem (7) of minimal norm in $L^2(\omega)$.

3.2 A Least-Squares Formulation of Problem (7)

A “reasonable” least-squares formulation of (7) reads as follows:

$$\begin{aligned} \text{Find } u \in L^2(\omega) \text{ such that} \\ J(u) \leq J(v), \quad \forall v \in L^2(\omega), \end{aligned} \quad (8)$$

with

$$J(v) = \frac{1}{2} \int_{\omega} [\alpha |\psi_2 - \psi_1|^2 + \mu |\nabla(\psi_2 - \psi_1)|^2] dx, \quad (9)$$

ψ_1 and ψ_2 in (9) being obtained from v via the solution of the elliptic boundary value problems (4) and (5), respectively. The functional J is clearly convex and C^∞ over $L^2(\omega)$. Any solution of problem (7) is a solution of the minimization problem (of the *virtual control* type) (8). Such a solution is characterized by

$$DJ(u) = 0, \quad (10)$$

where $DJ(\cdot)$ is the differential of functional J . Using classical methods from Control Theory (see, e.g., [2]), we can show that

$$\forall v \in L^2(\omega), DJ(v) = (p_1 - \psi_1)|_{\omega} + \psi_2, \quad (11)$$

where in (11), ψ_1 and ψ_2 are defined from v via the solution of (4) and (5), respectively, and where p_1 is the the unique solution of the following adjunct equation (written directly in variational form, here):

$$\begin{aligned} p_1 \in H_0^1(\Omega), \\ \int_{\Omega} [\alpha p_1 \varphi + \mu \nabla p_1 \cdot \nabla \varphi] dx = \int_{\omega} [\alpha (\psi_1 - \psi_2) \varphi + \mu \nabla (\psi_1 - \psi_2) \cdot \nabla \varphi] dx \\ + \frac{\mu}{l_s} \int_{\gamma} (\psi_2 - \psi_1) \varphi d\gamma, \quad \forall \varphi \in H_0^1(\Omega). \end{aligned} \quad (12)$$

4 On the Conjugate Gradient Solution of the Least-Squares Problem (8)

In order to solve the (linear) least-squares problem (8), we advocate a conjugate gradient algorithm operating in the space $L^2(\omega)$; this algorithm reads as follows:

$$u^0 \text{ is given in } L^2(\omega); \quad (13)$$

solve

$$\begin{aligned} \psi_1^0 &\in H^1(\Omega), \\ \alpha\psi_1^0 - \mu\nabla^2\psi_1^0 &= \tilde{f}(u^0) \text{ in } \Omega, \quad \psi_1^0 = g_0 \text{ on } \Gamma, \end{aligned} \quad (14)$$

$$\begin{aligned} \psi_2^0 &\in H^1(\omega), \\ \alpha\psi_2^0 - \mu\nabla^2\psi_2^0 &= u^0 \text{ in } \omega, \quad \mu\frac{\partial\psi_2^0}{\partial n} = \frac{\mu}{l_s}\psi_1^0 - g_1 \text{ on } \gamma, \end{aligned} \quad (15)$$

$$\begin{aligned} p_1^0 &\in H_0^1(\Omega), \\ \int_{\Omega} [\alpha p_1^0 \varphi + \mu \nabla p_1^0 \cdot \nabla \varphi] dx &= \int_{\omega} [\alpha(\psi_1^0 - \psi_2^0)\varphi + \mu \nabla(\psi_1^0 - \psi_2^0) \cdot \nabla \varphi] dx \\ &+ \frac{\mu}{l_s} \int_{\gamma} (\psi_2^0 - \psi_1^0)\varphi d\gamma, \quad \forall \varphi \in H_0^1(\Omega), \end{aligned} \quad (16)$$

and set

$$g^0 = (p_1^0 - \psi_1^0)|_{\omega} + \psi_2^0, \quad w^0 = g^0. \quad (17)$$

For $n \geq 0$, assuming that u^n, g^n and w^n are known, the last two different from 0, we compute u^{n+1}, g^{n+1} and w^{n+1} as follows (with χ_{ω} the characteristic function of ω):

Solve

$$\begin{aligned} \bar{\psi}_1^n &\in H_0^1(\Omega), \\ \alpha\bar{\psi}_1^n - \mu\nabla^2\bar{\psi}_1^n &= w^n \chi_{\omega} \text{ in } \Omega, \quad \bar{\psi}_1^n = 0 \text{ on } \Gamma, \end{aligned} \quad (18)$$

$$\begin{aligned} \bar{\psi}_2^n &\in H^1(\omega), \\ \alpha\bar{\psi}_2^n - \mu\nabla^2\bar{\psi}_2^n &= w^n \text{ in } \omega, \quad \mu\frac{\partial\bar{\psi}_2^n}{\partial n} = \frac{\mu}{l_s}\bar{\psi}_1^n \text{ on } \gamma, \end{aligned} \quad (19)$$

$$\begin{aligned} \bar{p}_1^n &\in H_0^1(\Omega), \\ \int_{\Omega} [\alpha\bar{p}_1^n \varphi + \mu \nabla\bar{p}_1^n \cdot \nabla \varphi] dx &= \int_{\omega} [\alpha(\bar{\psi}_1^n - \bar{\psi}_2^n)\varphi + \mu \nabla(\bar{\psi}_1^n - \bar{\psi}_2^n) \cdot \nabla \varphi] dx \\ &+ \frac{\mu}{l_s} \int_{\gamma} (\bar{\psi}_2^n - \bar{\psi}_1^n)\varphi d\gamma, \quad \forall \varphi \in H_0^1(\Omega), \end{aligned} \quad (20)$$

and set

$$\bar{g}^n = (\bar{p}_1^n - \bar{\psi}_1^n)|_{\omega} + \bar{\psi}_2^n. \quad (21)$$

Next, compute

$$\rho_n = \frac{\int_{\omega} |g^n|^2 dx}{\int_{\omega} \bar{g}^n w^n dx} \quad (22)$$

and

$$u^{n+1} = u^n - \rho_n w^n, \quad (23)$$

$$g^{n+1} = g^n - \rho_n \bar{g}^n. \quad (24)$$

If $\frac{\int_{\omega} |g^{n+1}|^2 dx}{\max\{\int_{\omega} |g^0|^2 dx, \int_{\omega} |u^{n+1}|^2 dx\}} \leq tol$, take $u = u^{n+1}$ and $\psi = \psi_1^{n+1}|_{\Omega \setminus \bar{\omega}}$; else compute

$$\gamma_n = \frac{\int_{\omega} |g^{n+1}|^2 dx}{\int_{\omega} |g^n|^2 dx} \quad (25)$$

and set

$$w^{n+1} = g^{n+1} + \gamma_n w^n. \quad (26)$$

Do $n + 1 \rightarrow n$ and return to (18).

5 On the Finite Element Implementation of the Least-Squares/ Fictitious Domain Methodology

5.1 Generalities

We (briefly) discuss in this section the finite element implementation of the least-squares/fictitious domain methodology described in Sects. 3 and 4. We will assume that $\bar{\omega} \subset \Omega \subset \mathbf{R}^2$ and that Ω is convex and/or has a smooth boundary. We assume also that γ is smooth. For simplicity, we still denote by ω and Ω the polygonal approximations of the above domains. From the triangulations \mathcal{T}_{h_1} of Ω and \mathcal{T}_{h_2} of ω , we define the following finite dimensional spaces:

$$V_{h_1} = \{\varphi | \varphi \in C^0(\bar{\Omega}), \varphi|_T \in P_1, \forall T \in \mathcal{T}_{h_1}\}, \quad (27)$$

$$V_{0h_1} = \{\varphi | \varphi \in V_{h_1}, \varphi = 0 \text{ on } \Gamma\}, \quad (28)$$

and

$$V_{h_2} = \{\varphi | \varphi \in C^0(\bar{\omega}), \varphi|_T \in P_1, \forall T \in \mathcal{T}_{h_2}\}, \quad (29)$$

P_1 being the space of the polynomials of two variables of degree ≤ 1 . We will use \mathbf{h} to denote the pair $\{h_1, h_2\}$. The finite element spaces V_{h_1} , V_{0h_1} and V_{h_2} are finite dimensional approximations to $H^1(\Omega)$, $H_0^1(\Omega)$ and $H^1(\omega)$, respectively. Similarly, we will use V_{h_2} to approximate the ‘‘control’’ space $L^2(\omega)$. As usual, h_1 (resp., h_2) denotes the length of the longest edge(s) of \mathcal{T}_{h_1} (resp., \mathcal{T}_{h_2}).

5.2 Finite Element Approximation of the Least-Squares Problem (8)

We approximate the least-squares problem (8) by

$$\begin{aligned} & \text{Find } u_{\mathbf{h}} \in V_{h_2} \text{ such that} \\ & J_{\mathbf{h}}(u_{\mathbf{h}}) \leq J_{\mathbf{h}}(v), \quad \forall v \in V_{h_2}, \end{aligned} \quad (30)$$

where

$$J_{\mathbf{h}}(v) = \frac{1}{2} \int_{\omega} [\alpha |\psi_2 - \pi_2 \psi_1|^2 + \mu |\nabla(\psi_2 - \pi_2 \psi_1)|^2] dx. \quad (31)$$

In (31), ψ_1 is the solution of the following fully discrete Dirichlet problem:

$$\begin{aligned} & \psi_1 \in V_{h_1}, \quad \psi_1 = g_{0h_1} \quad \text{on } \Gamma, \\ & \int_{\Omega} [\alpha \psi_1 \varphi + \mu \nabla \psi_1 \cdot \nabla \varphi] dx = \int_{\Omega} f_{h_1} \varphi dx + \int_{\omega} v \pi_2 \varphi dx, \quad \forall \varphi \in V_{0h_1}, \end{aligned} \quad (32)$$

where: (i) g_{0h_1} is an approximation of g_0 . (ii) $f_{h_1} \in V_{h_1}$; it approximates f over $\Omega \setminus \bar{\omega}$ and vanishes over ω . (iii) $\pi_2 : C^0(\bar{\Omega}) \rightarrow V_{h_2}$ is the interpolation operator defined as follows

$$\pi_2 \varphi = \sum_{i=1}^{N_{h_2}} \varphi(Q_i) w_{2i}, \quad \forall \varphi \in C^0(\bar{\Omega}), \quad (33)$$

$\{Q_i\}_{i=1}^{N_{h_2}}$ being the set of the vertices of \mathcal{T}_{h_2} and w_{2i} the P_1 -shape function associated with the vertex Q_i (we clearly have $N_{h_2} = \text{dimension of } V_{h_2}$). Returning to (31), the function ψ_2 there is the solution of the following discrete Neumann problem

$$\begin{aligned} & \psi_2 \in V_{h_2}, \\ & \int_{\omega} [\alpha \psi_2 \varphi + \mu \nabla \psi_2 \cdot \nabla \varphi] dx = \int_{\omega} v \varphi dx \\ & + \frac{\mu}{l_s} \int_{\gamma} (\pi_2 \psi_1 - g_{1h_2}) \varphi d\gamma, \quad \forall \varphi \in V_{h_2}, \end{aligned} \quad (34)$$

g_{1h_2} being an approximation of g_1 . As a discrete analogue of problem (8), the finite dimensional least-squares problem (30) is also of the virtual control type and well-suited to solution by a conjugate gradient algorithm operating in the space V_{h_2} . Due to page limitation we cannot describe this algorithm here; actually, this discrete analogue of algorithm (13)–(26) will be fully described in [1].

6 Numerical Experiments

As test problem, we consider the particular case of problem (1) associated with: (i) $\Omega = (0, 4) \times (0, 4)$, $\omega = \left\{ \{x_1, x_2\} \mid \left(\frac{x_1 - G_1}{a}\right)^2 + \left(\frac{x_2 - G_2}{b}\right)^2 < 1 \right\}$ with $G_1 = G_2 = 2$, $a = 1/4$ and $b = 1/8$. (ii) $f(x_1, x_2) = \alpha(x_1^3 - x_2^3) - 6\mu(x_1 - x_2)$, $\forall \{x_1, x_2\} \in \Omega \setminus \bar{\omega}$. (iii) $g_0 = x_1^3 - x_2^3$, $g_1 = \mu [3(n_1 x_1^2 - n_2 x_2^2) + (x_1^3 - x_2^3)/l_s]$, $\{n_1, n_2\} = \mathbf{n}$ being the unit normal vector at γ pointing to ω . (iv) $\alpha = 100$, $\mu = 0.1$

and $l_s = 0.1$. The unique solution of problem (1), associated with the above data, in $H^1(\Omega \setminus \overline{\omega})$ is given by

$$\psi(x_1, x_2) = x_1^3 - x_2^3.$$

Concerning the finite element implementation of the least-squares/fictitious domain methodology discussed in Sects. 3 and 4, we employed for \mathcal{T}_{h_1} (resp., \mathcal{T}_{h_2}) uniform triangulations of Ω (resp., triangulations of ω) like the one shown in Fig. 2(left) (resp., Fig. 2(right)). We used $u^0 = 0$ to initialize the discrete analogue of the conjugate gradient algorithm (13)–(26), and took $tol = 10^{-10}$ for the stopping criterion.

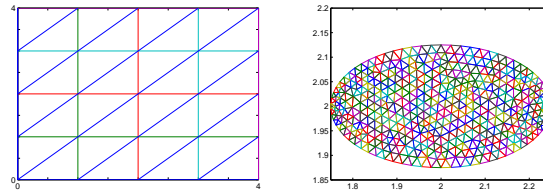


Fig. 2. A uniform triangulation of Ω (left); A triangulation of ω (right).

In Table 1, we report on the mesh sizes $h_2 = 1/40$ and $h_1 = 1/5, 1/10, 1/20$ and $1/40$: (i) The number of conjugate gradient iterations necessary to achieve convergence. (ii) Various norms of the approximation error. These results suggest: (a) For h_1 small enough, the number of iterations varies slowly with h_1 . (b) $\|\psi_{\mathbf{h}} - \psi\|_{L^\infty(\Omega \setminus \overline{\omega})} \approx O(h_1^2)$, $\|\psi_{\mathbf{h}} - \psi\|_{L^2(\Omega \setminus \overline{\omega})} \approx O(h_1^2)$ and $\|\psi_{\mathbf{h}} - \psi\|_{H^1(\Omega \setminus \overline{\omega})} \approx O(h_1)$. Concerning the decay of the cost function $J_{\mathbf{h}}$, we have, if $\mathbf{h} = \{1/10, 1/40\}$ (resp., $\mathbf{h} = \{1/20, 1/40\}$), $J_{\mathbf{h}}(u^0) = 3.67$ (resp., $J_{\mathbf{h}}(u^0) = 3.35$) and $J_{\mathbf{h}}(u^{61}) = 4.20 \times 10^{-9}$ (resp., $J_{\mathbf{h}}(u^{59}) = 5.11 \times 10^{-9}$), showing clearly that the computed approximations of ψ_1 and ψ_2 match quite well over ω . In order to further investigate the convergence properties of the methodology discussed in the above sections we performed computations with $h_2 = 1/20$ and $h_1 = 1/10, 1/20, 1/40$ and $1/80$. The corresponding results have been reported in Table 2. From these results we observe that: (i) If $h_1 \geq h_2$, the number of iterations necessary to achieve convergence does not vary significantly with h_1 ; on the other hand this number of iterations increases sharply when h_1 decreases below h_2 . (ii) The various approximation errors vary as expected (that is as in Table 1) if $h_1 \geq h_2$; on the other hand, they vary quite differently if $h_1 < h_2$, the only one behaving “nicely” being $\|\psi_{\mathbf{h}} - \psi\|_{H^1(\Omega \setminus \overline{\omega})}$, which shows a text-book $O(h_1)$ behavior as h_1 varies over the interval $[1/80, 1/10]$. From these results we suggest to take $h_1 = h_2$ to be on the safe side.

Remark 2. The results reported in [1] show a sharp decrease of the number of iteration when the methodology discussed here is applied to the solution of parabolic problems, including situations where ω is moving.

Table 1. Summary of numerical results ($h_2 = 1/40$).

h_1	Number of iterations	$\ \psi - \psi_h\ _{L^\infty(\Omega \setminus \overline{\omega})}$	$\ \psi - \psi_h\ _{L^2(\Omega \setminus \overline{\omega})}$	$\ \psi - \psi_h\ _{H^1(\Omega \setminus \overline{\omega})}$
1/5	34	0.1046	7.8370E-03	0.2855
1/10	61	2.1845E-02	1.9028E-03	0.1423
1/20	59	4.5840E-03	4.7015E-04	7.1089E-02
1/40	68	1.1385E-03	1.1708E-04	3.5518E-02

Table 2. Summary of numerical results ($h_2 = 1/20$).

h_1	Number of iterations	$\ \psi - \psi_h\ _{L^\infty(\Omega \setminus \overline{\omega})}$	$\ \psi - \psi_h\ _{L^2(\Omega \setminus \overline{\omega})}$	$\ \psi - \psi_h\ _{H^1(\Omega \setminus \overline{\omega})}$
1/10	33	2.1845E-02	1.9038E-03	0.1424
1/20	36	4.5840E-03	4.6807E-04	7.1063E-02
1/40	114	2.1385E-03	1.0163E-04	3.5434E-02
1/80	85	3.1514E-03	5.2854E-05	1.7532E-02

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