

# Multigrid Interpretation of a Three-Level Parareal Algorithm

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## 1 Introduction

Parallel-in-time methods, of which parareal [13] and multigrid reduction in time (MGRIT) [3] are well-known examples, are important tools for increasing parallelism beyond traditional spatially parallel methods, see [6, 14] and references therein. As a two-level method, parareal performs the fine but expensive integration independently (and in parallel) over many short time intervals, and it uses a cheap (but coarse) integrator to correct values across time subintervals sequentially. For linear ODE systems, parareal iterates are known to be equivalent to two-level MGRIT ones for a specific choice of initial guess, restriction/prolongation operators and relaxation scheme, cf. [3, 9, 10]. One can thus analyze parareal convergence in two ways: one can make hypotheses on Lipschitz constants and truncation errors, which is typical in the ODE community, cf. [1, 8, 13], or one can use spectral information of all-at-once matrices, as is common in the multigrid community, see [2, 3, 5, 15].

When parareal and MGRIT are used with many time subintervals, the coarse correction step becomes a computational bottleneck. To overcome this, one can parallelize the coarse solution by subdividing the coarse problem and using a coarser level to ensure global communication. For MGRIT, this leads to a multilevel variant [11]; for parareal, a three-level variant has been introduced and analyzed in [12]. In this paper, we show that there is a choice of restriction/prolongation operators and relaxation schemes such that the resulting MGRIT method is equivalent to three-level parareal when applied to linear problems. The existing MGRIT literature can thus add to our understanding of three-level parareal, beyond what is shown in [12].

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## 2 The three-level parareal algorithm

Suppose one wishes to solve the *linear* system of ODEs  $u' = \Phi u + f(t)$  with initial conditions  $u(0) = u_0$  on the interval  $[0, T]$ . To obtain the temporal grid for both parareal and MGRIT, we subdivide the interval hierarchically as follows:<sup>1</sup>

- The interval  $[0, T]$  is subdivided into  $p$  coarsest intervals  $\mathcal{I}_i = [T_{i-1}, T_i]$ ,  $i = 1, \dots, p$ , each of length  $\Delta T = T/p$ ;
- Each coarsest interval  $\mathcal{I}_i$  is subdivided into  $m$  subintervals  $\mathcal{I}_{i,j} = [t_{i,j}, t_{i,j+1}]$ ,  $j = 0, 1, \dots, m-1$ , of length  $\Delta t = \Delta T/m$ ;
- Each  $\mathcal{I}_{i,j}$  is divided into intervals  $[t_{i,j,k}, t_{i,j,k+1}]$  ( $0 \leq k < n$ ) of length  $\delta t = \Delta t/n$ .

We can now define the following *propagators*, which take an initial value at the beginning of  $\mathcal{I}_i$ ,  $\mathcal{I}_{i,j}$  or  $\mathcal{I}_{i,j,k}$  and return the solution at the end of the interval:<sup>2</sup>

- $F_0$  is the action of the fine integrator over one fine time step  $\delta t$ . For a linear problem, we have  $F_0 u_{i-1} = \Phi_0 u_{i-1} + f_i$ .
- $F = F_0^n$  is the action of the fine integrator over one intermediate time step  $\Delta t = n\delta t$ . For a linear problem, we have  $F u_{i-n} = \Phi_0^n u_{i-n} + \sum_{k=0}^{n-1} \Phi_0^k f_{i-k}$ .
- $G$  is the action of the intermediate integrator over one intermediate time step  $\Delta t$ . For a linear problem, we have  $G U_{i,j-1} = \Phi_1 U_{i,j-1} + \gamma_{i,j}$ .
- $H$  is the action of the coarse integrator over one coarse time step  $\Delta T = m\Delta t$ . For a linear problem, we have  $H Y_{i-1} = \Phi_2 Y_{i-1} + \eta_i$ .

The three-level parareal algorithm, as introduced in [12], iterates on the level-1 state variables  $U_{i,j}$  and level-2 state variables  $Y_i$  as follows:

1. Initialization (with iteration indices appearing as superscripts):

$$\begin{aligned} Y_0^0 &= u_0, & Y_i^0 &= H Y_{i-1}^0 \\ U_{i,0}^0 &= Y_{i-1}^0, & U_{i,j}^0 &= G U_{i,j-1}^0 \quad (1 \leq j \leq m), \end{aligned}$$

2. Iteration: for  $\nu = 0, 1, 2, \dots$ ,

$$U_{i,0}^{\nu+1} = Y_{i-1}^{\nu}, \quad U_{i,j}^{\nu+1} = F U_{i,j-1}^{\nu} + G U_{i,j-1}^{\nu+1} - G U_{i,j-1}^{\nu} \quad (1 \leq j \leq m), \quad (1)$$

$$Y_i^{\nu+1} = u_0, \quad Y_i^{\nu+1} = U_{i,m}^{\nu+1} + H Y_{i-1}^{\nu+1} - H Y_{i-1}^{\nu}. \quad (2)$$

This method is shown in [12] to converge to the fine solution *in finitely many steps*, i.e.,  $U_{i,j}^{\nu} = F^{(i-1)m+j} u_0$  for  $\nu \geq i(m+1)$ , for any choice of  $G$  and  $H$ . Note that this is not a nested iteration, where one needs to iterate  $U$  or  $Y$  to sufficient accuracy before switching levels; instead, only one parareal step on  $U_{i,j}$  is performed before it is used in (2), and one coarse parareal step (2) is performed before the  $Y_i$  are used as new initial values in (1).

<sup>1</sup> For ease of explanation, we assume that all subdivisions have equal length, although it is easy to see that similar results hold for non-uniform subdivisions.

<sup>2</sup> To lighten the notation, the time index is only indicated in the variable on which the propagators are applied, and not in the propagators themselves.

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**Algorithm 1** MGRIT( $\ell, \tilde{\mathbf{g}}^{(\ell)}$ ) (in correction form, as defined in [3])

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**if**  $\ell$  is the coarsest level  $L$  **then**  
  Solve coarse grid system  $A_L \mathbf{u}^{(L)} = \tilde{\mathbf{g}}^{(L)}$   
**else**  
  Relax on  $A_\ell \mathbf{u}^{(\ell)} = \tilde{\mathbf{g}}^{(\ell)}$  using  $F$ -relaxation  
  Compute and restrict residual using injection:  $\tilde{\mathbf{g}}^{(\ell+1)} = R_\ell^{\ell+1}(\tilde{\mathbf{g}}^{(\ell)} - A_\ell \mathbf{u}^{(\ell)})$   
  Solve on the next level : MGRIT( $\ell + 1, \tilde{\mathbf{g}}^{(\ell+1)}$ )  
  Correct:  $\mathbf{u}^{(\ell)} \leftarrow \mathbf{u}^{(\ell)} + P_{\ell+1}^\ell \mathbf{u}^{(\ell+1)}$   
**end if**

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**Algorithm 2** MGRIT-FAS( $\ell, \mathbf{u}^{(\ell)}, \mathbf{g}^{(\ell)}$ ) (as defined in [4])

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**if**  $\ell$  is the coarsest level  $L$  **then**  
  Solve coarse grid system  $A_L(\mathbf{u}^{(L)}) = \mathbf{g}^{(L)}$   
**else**  
  Relax on  $A_\ell(\mathbf{u}^{(\ell)}) = \mathbf{g}^{(\ell)}$  using  $F$ -relaxation to obtain  $\mathbf{v}^{(\ell)}$   
  Compute FAS right hand side:  $\mathbf{g}^{(\ell+1)} = R_\ell^{\ell+1}(\mathbf{g}^{(\ell)} - A_\ell(\mathbf{v}^{(\ell)})) + A_{\ell+1}(R_\ell^{\ell+1} \mathbf{v}^{(\ell)})$   
  Solve on the next level : MGRIT-FAS( $\ell + 1, \mathbf{u}^{(\ell+1)}, \mathbf{g}^{(\ell+1)}$ )  
  Correct:  $\mathbf{u}^{(\ell)} \leftarrow \mathbf{u}^{(\ell)} + P_{\ell+1}^\ell(\mathbf{u}^{(\ell+1)} - R_\ell^{\ell+1} \mathbf{v}^{(\ell)})$   
**end if**

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### 3 Equivalence with the MGRIT V-cycle

The initial value problems that are solved by the propagators can also be written as linear systems of the type  $A_\ell \mathbf{u}^{(\ell)} = \mathbf{g}^{(\ell)}$ , where

$$A_\ell = \begin{bmatrix} I & & & & \\ -\Phi_\ell & I & & & \\ & & \ddots & \ddots & \\ & & & & -\Phi_\ell & I \end{bmatrix}.$$

The index  $\ell$  here indicates the level of coarseness of the temporal grid, with  $\ell = 0$  being the finest grid, and  $\ell = 2$  being the coarsest for a three-level method. Such systems can be solved using the MGRIT V-cycle with F-relaxation algorithm, which can be written in correction form [3] or as a full approximation scheme (FAS) [4], see Algorithms 1 and 2. Here, we consider the special case of  $L = 2$ , i.e., the three-level algorithm. For the purpose of writing the recurrence, we will index the fine grid (level-0) solution as  $u_{i,j,k} \approx u(t_{i,j,k})$ . The level-1 vectors will be double indexed as  $u_{i,j} \approx u(t_{i,j})$ , and level-2 vectors are singly indexed as  $u_i \approx u(T_{i-1})$ . If injection is used for  $P_{\ell+1}^\ell$  and  $R_\ell^{\ell+1} = (P_{\ell+1}^\ell)^T$  in Algorithm 2, then one V-cycle of MGRIT-FAS with F-relaxation for solving  $A_0 \mathbf{u} = f$  updates the iterate  $u_{i,j,k}$  as follows:

1. Relax on level 0:

$$v_{i,j,k} = \begin{cases} f_{i,j,k} + \Phi_0 v_{i,j,k-1}, & 1 \leq k \leq n-1, \quad \forall i, j, \\ u_{i,j,0}, & k = 0. \end{cases}$$

2. Compute FAS right-hand side for level 1:

$$g_{i,j} = \begin{cases} f_{i,j,0} + \Phi_0 v_{i,j-1,n-1} - \Phi_1 u_{i,j-1,0}, & 1 \leq j \leq m-1, \\ f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} - \Phi_1 u_{i-1,m-1,0}, & j = 0. \end{cases}$$

3. Relax on level 1 using initial guess  $(\mathbf{u}^{(1)})_{i,j} = v_{i,j,0} = u_{i,j,0}$ :

$$v_{i,j} = \begin{cases} f_{i,j,0} + \Phi_0 v_{i,j-1,n-1} + \Phi_1 (v_{i,j-1} - u_{i,j-1,0}), & 1 \leq j \leq m-1, \\ u_{i,0,0}, & j = 0. \end{cases}$$

4. Compute FAS right-hand side for level 2:

$$g_i = f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} + \Phi_1 (v_{i-1,m-1} - u_{i-1,m-1,0}) - \Phi_2 u_{i-1,0,0}.$$

5. Solve the level-2 system:

$$u_i^{\text{new}} = f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} + \Phi_1 (v_{i-1,m-1} - u_{i-1,m-1,0}) + \Phi_2 (u_{i-1}^{\text{new}} - u_{i-1,0,0}).$$

6. Correct on level 1 and then on level 0, using injection for both levels: we set for all  $1 \leq i \leq p$

$$u_{i,j,k}^{\text{new}} = \begin{cases} f_{i,j,k} + \Phi_0 v_{i,j,k-1}, & 1 \leq k \leq n-1, \forall j, \\ f_{i,j,0} + \Phi_0 v_{i,j-1,n-1} + \Phi_1 (v_{i,j-1} - u_{i,j-1,0}), & k = 0, 1 \leq j \leq m-1, \\ u_{i,0,0}^{\text{new}} = f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} + \Phi_1 (v_{i-1,m-1} - u_{i-1,m-1,0}) + \Phi_2 (u_{i-1}^{\text{new}} - u_{i-1,0,0}). \end{cases}$$

We can now prove the following equivalence theorem.

**Theorem 1** For the linear problem  $u' = \Phi u + f(t)$ , assume that  $u_{i,j,k}^0$  satisfies

$$u_{1,0,0}^0 = u_0, \quad u_{i,0,0}^0 = H u_{i-1,0,0}^0 \quad \forall i \geq 1, \quad u_{i,j,0}^0 = G u_{i,j-1,0}^0 \quad \forall j = 1, \dots, m-1.$$

Then for all  $\nu \geq 0$ , the three-level MGRIT-FAS V-cycle with F-relaxation and with injection as the prolongation operator is equivalent to three-level parareal via

$$u_{i,j,k}^{\nu+1} = \begin{cases} F_0^k U_{i,j}^\nu, & 1 \leq k \leq n-1, \quad \forall i, j, \\ U_{i,j}^{\nu+1}, & k = 0, 1 \leq j \leq m-1, \quad \forall i \geq 1, \\ Y_{i-1}^{\nu+1}, & j = k = 0, \quad \forall i \geq 1. \end{cases}$$

**Proof** From the initialization conditions, we have for  $\nu = 0$  that  $u_{i,j,0}^\nu = U_{i,j}^\nu$  for  $1 \leq j \leq m-1$ , and  $u_{i,0,0}^\nu = Y_{i-1}^\nu$  for all  $i$ . We will prove by induction that these two equalities also hold for  $\nu \geq 1$ . To do so, we rewrite  $u_{i,j,k}^{\text{new}}$  in terms of the propagators  $F_0, F, G$  and  $H$ . The update formula at step 6 leads us to consider three cases:

**Case 1 ( $k \neq 0$ ).** Step 1 at iteration  $\nu$  reads

$$u_{i,j,k}^{\text{new}} = v_{i,j,k} = F_0 v_{i,j,k-1} = \dots = F_0^k v_{i,j,0} = F_0^k U_{i,j}^\nu.$$

**Case 2 ( $k = 0, j \neq 0$ ).** This case is given by step 3, where

$$u_{i,j,0}^{\text{new}} = v_{i,j} = F_0 v_{i,j-1,n-1} + \Phi_1(v_{i,j-1} - U_{i,j-1}^Y) = F_0^n U_{i,j-1}^Y + G v_{i,j-1} - G U_{i,j-1}^Y.$$

Here, we have replaced the difference of  $\Phi_1$  by a difference of  $G$ , because  $G$  is affine. Thus, we have  $v_{i,j} = U_{i,j}^{Y^{v+1}}$  for  $1 \leq j \leq m-1$ , since both quantities are initialized the same way (we have  $v_{i,0} = u_{i,0,0} = Y_{i-1}^Y = U_{i,0}^{Y^{v+1}}$ ) and satisfy the same recurrence.

**Case 3 ( $j = k = 0$ ).** Here we have  $u_{i,0,0}^{\text{new}} = u_i^{\text{new}}$ , so step 5 gives, for  $i \geq 2$ ,

$$\begin{aligned} u_i^{\text{new}} &= f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} + \Phi_1(v_{i-1,m-1} - u_{i-1,m-1,0}) + \Phi_2(u_{i-1}^{\text{new}} - u_{i-1,0,0}) \\ &= F_0 v_{i-1,m-1,n-1} + \Phi_1(v_{i-1,m-1} - U_{i-1,m-1}^Y) + \Phi_2(u_{i-1}^{\text{new}} - u_{i-1,0,0}) \\ &= F_0^n U_{i-1,m-1}^Y + G v_{i-1,m-1} - G U_{i-1,m-1}^Y + H u_{i-1}^{\text{new}} - H u_{i-1,0,0} \\ &= U_{i-1,m}^{Y^{v+1}} + H u_{i-1}^{\text{new}} - H Y_{i-2}^Y. \end{aligned}$$

For  $i = 1$ , we have  $u_1^{\text{new}} = u_0 = Y_0^{Y^{v+1}}$ ; thus,  $u_i^{\text{new}}$  and  $Y_{i-1}^{Y^{v+1}}$  satisfy the same recurrence with the same initial condition. This leads to  $u_{i,0,0}^{\text{new}} = Y_{i-1}^{Y^{v+1}}$  for all  $i$ , as claimed.  $\square$

We can now use the FAS formulation to deduce the equivalence in classical (correction) form. We define the following operators:

$$E_\ell = I - P_{\ell+1}^\ell R_\ell^{\ell+1}, \quad M_\ell = \text{diag}((A_\ell)_{11}, (A_\ell)_{22}, \dots),$$

where  $(A_\ell)_{ii}$  are diagonal blocks of  $A_\ell$  corresponding to the  $i$ th subinterval, *starting* with the coarse point and including all the fine points until (but excluding) the next coarse point. In other words,  $M_\ell$  is the block Jacobi smoother for level  $\ell$ , and  $E_\ell$  blanks out the coarse points and retains the fine points when applied to a vector of values at level  $\ell$ . Similar operators were defined in [10], where the authors proved the equivalence between two-level parareal and a geometric multigrid method with block Jacobi smoothing and aggressive coarsening in the FAS setting; however, the blocks in [10] are defined differently, with the coarse points appearing at the end of the block rather than the beginning. We write the change in the solution at step 6 as

$$u_{i,j,k}^{\text{new}} - u_{i,j,k} = \begin{cases} v_{i,j,k} - u_{i,j,k} =: (\Delta \mathbf{u}^{(0)})_{i,j,k}, & k \neq 0, \\ v_{i,j} - u_{i,j,0} =: (\Delta \mathbf{u}^{(1)})_{i,j}, & k = 0, j \neq 0, \\ u_i^{\text{new}} - u_{i,0,0} =: (\Delta \mathbf{u}^{(2)})_i, & j = k = 0. \end{cases}$$

To compute  $\Delta \mathbf{u}^{(0)}$ , note that  $v_{i,j,k} - u_{i,j,k} = 0$  when  $k = 0$ ; for  $k \neq 0$ , we have

$$\begin{aligned} (\Delta \mathbf{u}^{(0)})_{i,j,k} &= v_{i,j,k} - u_{i,j,k} = f_{i,j,k} + \Phi_0(v_{i,j,k-1} - u_{i,j,k-1}) + \Phi_0 u_{i,j,k-1} - u_{i,j,k} \\ &= (\mathbf{f} - A_0 \mathbf{u})_{i,j,k} + \Phi_0 (\Delta \mathbf{u}^{(0)})_{i,j,k-1}. \end{aligned}$$

If we move  $\Phi_0 (\Delta \mathbf{u}^{(0)})_{i,j,k-1}$  to the left and recall the definition of  $M_0$ , we get

$$M_0 \Delta \mathbf{u}^{(0)} = E_0 (\mathbf{f} - A_0 \mathbf{u}) \implies \Delta \mathbf{u}^{(0)} = M_0^{-1} E_0 \tilde{\mathbf{g}}^{(0)},$$

where  $\tilde{\mathbf{g}}^{(0)} = \mathbf{f} - A_0\mathbf{u}$  is the initial residual. This is almost the same as in [10], except the residual is blanked before the smoothing, instead of after. Next, we calculate

$$(\Delta\mathbf{u}^{(1)})_{i,j} = v_{i,j} - u_{i,j,0} = \begin{cases} 0, & j = 0, \\ g_{i,j} + \Phi_1(v_{i,j-1} - u_{i,j-1,0}) + \Phi_1 u_{i,j-1,0} - u_{i,j,0}, & j \neq 0, \end{cases}$$

which implies

$$M_1\Delta\mathbf{u}^{(1)} = E_1(\mathbf{g}^{(1)} - A_1R_0^1\mathbf{u}) = E_1R_0^1(\mathbf{f}^{(0)} - A_0(\mathbf{u} + \Delta\mathbf{u}^{(0)})).$$

Thus,  $\Delta\mathbf{u}^{(1)} = M_1^{-1}E_1\tilde{\mathbf{g}}^{(1)}$ , where  $\tilde{\mathbf{g}}^{(1)} = R_0^1(\tilde{\mathbf{g}}^{(0)} - A_0\Delta\mathbf{u}^{(0)})$ . Finally, we have

$$(\Delta\mathbf{u}^{(2)})_i = u_i^{\text{new}} - u_{i,0,0} = g_i + \Phi_2(u_{i-1}^{\text{new}} - u_{i-1,0,0}) + \Phi_2 u_{i-1,0,0} - u_{i,0,0},$$

which leads to

$$A_2\Delta\mathbf{u}^{(2)} = \mathbf{g}^{(2)} - A_2R_0^2\mathbf{u} = R_1^2(\mathbf{g}^{(1)} - A_1(R_0^1\mathbf{u} + \Delta\mathbf{u}^{(1)})) = R_1^2(\tilde{\mathbf{g}}^{(1)} - A_1\Delta\mathbf{u}^{(1)}).$$

We conclude, by replacing  $\Delta\mathbf{u}^{(1)}$  with  $M_1^{-1}E_1\tilde{\mathbf{g}}^{(1)}$  in the last step, that

$$\begin{aligned} \mathbf{u}^{\text{new}} - \mathbf{u} &= \Delta\mathbf{u}^{(0)} + P_1^0\Delta\mathbf{u}^{(1)} + P_2^0\Delta\mathbf{u}^{(2)} \\ &= \Delta\mathbf{u}^{(0)} + P_1^0(\Delta\mathbf{u}^{(1)} + P_2^1A_2^{-1}R_1^2(\tilde{\mathbf{g}}^{(1)} - A_1\Delta\mathbf{u}^{(1)})) \\ &= \Delta\mathbf{u}^{(0)} + P_1^0((I - P_2^1A_2^{-1}R_1^2A_1)M_1^{-1}E_1 + P_2^1A_2^{-1}R_1^2)\tilde{\mathbf{g}}^{(1)} \end{aligned}$$

Defining  $T = (I - P_2^1A_2^{-1}R_1^2A_1)M_1^{-1}E_1 + P_2^1A_2^{-1}R_1^2$ , we continue to calculate

$$\begin{aligned} \mathbf{u}^{\text{new}} - \mathbf{u} &= \Delta\mathbf{u}^{(0)} + P_1^0TR_0^1(\tilde{\mathbf{g}}^{(0)} - A_0\Delta\mathbf{u}^{(0)}) \\ &= (P_1^0TR_0^1 + (I - P_1^0TR_0^1A_0)M_0^{-1}E_0)(\mathbf{f} - A\mathbf{u}) =: \mathcal{P}(\mathbf{f} - A\mathbf{u}). \end{aligned}$$

We conclude that the error propagator reads

$$S = I - \mathcal{P}A_0 = (I - P_1^0TR_0^1A_0)(I - M_0^{-1}E_0A_0),$$

where the operator  $T$  satisfies  $I - TA_1 = (I - P_2^1A_2^{-1}R_1^2A_1)(I - M_1^{-1}E_1A_1)$ . Note that the preconditioners  $\mathcal{P}$  and  $T$  can also be written as

$$\mathcal{P} = M_0^{-1}E_0 + P_1^0TR_0^1(I - A_0M_0^{-1}E_0), \quad T = M_1^{-1}E_1 + P_2^1A_2^{-1}R_1^2(I - A_1M_1^{-1}E_1).$$

We can hence interpret the action of the preconditioner  $\mathcal{P}$  as follows:

1.  $M_0^{-1}E_0$ : Take the fine residual, blank out the coarse points and apply block Jacobi.
2.  $I - A_0M_0^{-1}E_0$ : Update the residual after relaxation.
3.  $P_1^0TR_0^1$ : Restrict the new residual, recursively solve the coarse problem, then update the coarse points by injection.

Since  $T$  acts the same way but at a coarser level, the action of  $\mathcal{P}$  corresponds to exactly one MGRIT V-cycle with  $F$ -relaxation, written in correction form.

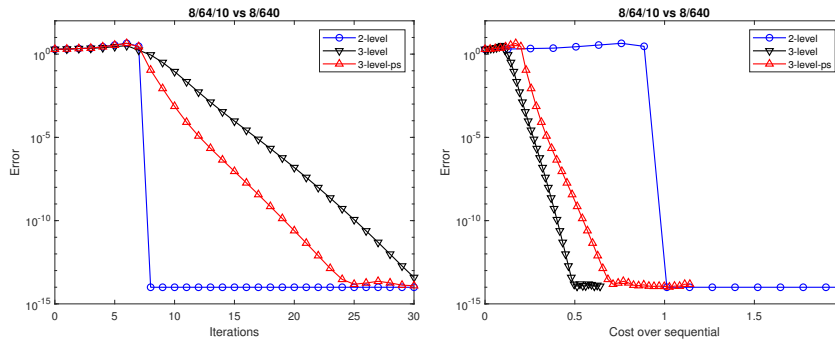
**Remark** If one replaces injection with injection plus  $F$ -relaxation (like in standard MGRIT), then the equivalent parareal formulation at the  $\nu$ th iteration would be

$$\begin{aligned} U_{i,0}^{\nu+1/2} &= Y_{i-1}^\nu, & U_{i,j}^{\nu+1/2} &= GU_{i,j-1}^{\nu+1/2} + FU_{i,j-1}^\nu - GU_{i,j-1}^\nu & (1 \leq j \leq m), \\ Y_0^{\nu+1} &= u_0, & Y_i^{\nu+1} &= U_{im}^{\nu+1/2} + HY_{i-1}^{\nu+1} - HY_{i-1}^\nu, \\ U_{i,0}^{\nu+1} &= Y_{i-1}^{\nu+1}, & U_{i,j}^{\nu+1} &= GU_{i,j-1}^{\nu+1} + FU_{i,j-1}^\nu - GU_{i,j-1}^\nu & (1 \leq j \leq m). \end{aligned}$$

Note that the term  $FU_{i,j-1}^\nu - GU_{i,j-1}^\nu$  is used twice, but it only needs to be computed once using a fine propagation. The intermediate propagation  $G$ , however, needs to be computed twice, since it is applied once to  $U_{i,j-1}^{\nu+1/2}$ , and another time to  $U_{i,j-1}^{\nu+1}$ .

## 4 Numerical example

We present the numerical example in [7], where the advection-diffusion equation  $u_t = u_x + \kappa u_{xx}$  with periodic boundary conditions  $u(0,t) = u(2,t)$ ,  $u_x(0,t) = u_x(2,t)$  is solved on  $t \in (0,4)$ , with  $\kappa = 1/1024$  (advection-dominated case) and  $u(x,0) = e^{-20(x-1)^2}$ . We discretize the problem using second order finite difference in space and backward Euler in time, with  $\Delta x = 1/20$  and  $\delta t = 1/1280$ . For two-level parareal, the coarse propagator is backward Euler with  $\Delta T = 1/2$  (8 coarse steps with 640 fine steps per coarse step). For three-level parareal, we use an intermediate level with  $\Delta t = 1/128$  (10 fine steps per intermediate step), while keeping  $\Delta T = 1/2$  for the coarsest level (i.e., 64 intermediate steps per coarse step). In Figure 1, we compare two-level and three-level parareal, both with and without post-smoothing. We compare both the iteration count and the *idealized* running time, as measured by the number of *non-concurrent* backward Euler steps taken at all levels; this cost is normalized by that of sequential time-stepping, so that a cost of 1 means the same cost as sequential time-stepping without parallelization. We see that two-level parareal



**Fig. 1** Left: Iteration count for two-level parareal, and three-level parareal, with and without post-smoothing. Right: Computational cost of the three methods, as measured by the number of backward Euler steps taken, normalized by the cost of sequential time-stepping.

converges to the exact solution in 8 iterations, whereas the three-level variants take many more iterations. However, the three-level iterations are much more parallel and take less time to run than a two-level iteration. In particular, both three-level versions converge with cost much lower than 1; such speedup is not possible for two-level parareal. Finally, although post-smoothing reduces the number of three-level iterations, the higher cost per iteration (two intermediate propagations rather than one) makes it slower than no post-smoothing once the normalized cost is considered.

**Acknowledgements** F. Kwok acknowledges support from the National Science and Engineering Research Council of Canada (RGPIN- 2021-02595). The work described in this paper is partially supported by a grant from the ANR/RGC joint research scheme sponsored by the Research Grants Council of the Hong Kong Special Administrative Region, China and the French National Research Agency (Project no. A-CityU203/19). This project has received funding from the Federal Ministry of Education and Research and the European High-Performance Computing Joint Undertaking (JU) under grant agreement No 955701, Time-X. The JU receives support from the European Union's Horizon 2020 research and innovation programme and Belgium, France, Germany, Switzerland.

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