# A Smoother Based on Nonoverlapping Domain Decomposition Methods for H(div) Problems: A Numerical Study

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**Abstract** The purpose of this paper is to introduce a V-cycle multigrid method for vector field problems discretized by the lowest order Raviart-Thomas hexahedral element. Our method is connected with a smoother based on a nonoverlapping domain decomposition method. We present numerical experiments to show the effectiveness of our method.

## **1** Introduction

Let  $\Omega$  be a bounded domain in  $\mathbb{R}^3$  and  $H_0(\operatorname{div}; \Omega)$  be the space of square integrable vector fields on  $\Omega$  that have square integrable divergence in  $\Omega$  and vanishing normal components on  $\partial \Omega$  (cf. [7]). In this paper we consider a multigrid method for the following problem: Find  $\boldsymbol{u} \in H_0(\operatorname{div}; \Omega)$  such that

$$a(\boldsymbol{u},\boldsymbol{v}) = (\boldsymbol{f},\boldsymbol{v}) \qquad \forall \boldsymbol{v} \in H_0(\operatorname{div};\Omega), \tag{1}$$

where

$$a(\boldsymbol{w},\boldsymbol{v}) = \boldsymbol{\alpha}(\operatorname{div}\boldsymbol{w},\operatorname{div}\boldsymbol{v}) + \boldsymbol{\beta}(\boldsymbol{w},\boldsymbol{v}), \qquad (2)$$

and  $(\cdot, \cdot)$  is the inner product on  $L_2(\Omega)$  (or  $[L_2(\Omega)]^3$ ). We assume that  $\mathbf{f} \in [L_2(\Omega)]^3$ and  $\alpha$  and  $\beta$  are positive. Unlike the scalar elliptic equation case, multigrid methods for the problem (1) with simple smoothers do not work. We need a special treatment

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for the smoother. In [2–4, 9], an overlapping domain decomposition preconditioner was employed in the construction of the smoother.

Our goal is to develop multigrid methods in the same spirit but using nonoverlapping domain decomposition preconditioners instead, which reduce the dimensions of the subproblems that have to be solved. We note that other multigrid methods for H(div) were investigated in [8, 10].

Applications of fast solvers for H(div) problems are discussed for example in [2, 11–13, 16]. In particular the multigrid method in this paper can be applied to a mixed method for second order partial differential equations based on a first-order system least-squares formulation [2, 6], which is equivalent to our model problem. It can also be used as an effective preconditioner for H(div) problems with variable coefficients. The model problem also arise in Reissner-Mindlin plates [1] and Brinkman equations [15].

In [5], there are similar ingredients and convergence analysis for the convex domain and the constant coefficient case. In this paper, we mainly focus on the numerical study that is not covered by the theory in [5].

The rest of this paper is organized as follows. We present the standard discretization of (1) by the lowest order Raviart-Thomas hexahedral element in Section 2. We next introduce the V-cycle multigrid method in Section 3. Finally, numerical experiments are presented in Section 4.

### 2 The Discrete Problem

Let  $\mathscr{T}_h$  be a hexahedral triangulation of  $\Omega$ . The lowest order Raviart-Thomas H(div) conforming finite element space [14] is denoted by  $V_h$ . A vector field  $\boldsymbol{v}$  belongs to  $V_h$  if and only if it belongs to  $H_0(\text{div};\Omega)$  and takes the form

$$\begin{bmatrix} a_1 \\ a_2 \\ a_3 \end{bmatrix} + \begin{bmatrix} b_1 x_1 \\ b_2 x_2 \\ b_3 x_3 \end{bmatrix}$$

on each hexahedral element, where the  $a_i$ 's and  $b_i$ 's are constants. On each hexahedral element T the vector field v is determined by the six degrees of freedom defined by the average of the normal component on each face. The discrete problem for (1) is to find  $u_h \in V_h$  such that

$$a(\boldsymbol{u}_h, \boldsymbol{v}) = \int_{\Omega} \boldsymbol{f} \cdot \boldsymbol{v} \, dx \qquad \forall \, \boldsymbol{v} \in V_h. \tag{3}$$

In the multigrid approach we solve (3) on a sequence of triangulations  $\mathcal{T}_0, \mathcal{T}_1, \ldots$ , where  $\mathcal{T}_0$  is an initial triangulation of  $\Omega$  by hexahedral elements and  $\mathcal{T}_k$   $(k \ge 1)$ is obtained from  $\mathcal{T}_{k-1}$  by uniform subdivision. We will denote the lowest order Raviart-Thomas finite element space associated with  $\mathcal{T}_k$  by  $V_k$ . The *k*-th level discrete problem is to find  $\boldsymbol{u}_k \in V_k$  such that A Multigrid for H(div)

$$a(\boldsymbol{u}_k, \boldsymbol{v}) = (\boldsymbol{f}, \boldsymbol{v}) \qquad \forall \, \boldsymbol{v} \in V_k.$$

Let  $A_k: V_k \longrightarrow V'_k$  be defined by

$$\langle A_k \boldsymbol{w}, \boldsymbol{v} \rangle = a(\boldsymbol{w}, \boldsymbol{v}) \qquad \forall \boldsymbol{v}, \boldsymbol{w} \in V_k, \tag{4}$$

where  $\langle \cdot, \cdot, \rangle$  is the canonical bilinear form on  $V'_k \times V_k$ . We can then rewrite the *k*-th level discrete problem as

$$A_k \boldsymbol{u}_k = f_k, \tag{5}$$

where  $f_k \in V'_k$  is defined by

$$\langle f_k, \boldsymbol{v} \rangle = (\boldsymbol{f}, \boldsymbol{v}) \qquad \forall \boldsymbol{v} \in V_k.$$

Multigrid methods are optimal order iterative methods for equations of the form

$$A_k \mathbf{z} = g \tag{6}$$

that includes (5) as a special case.

### 3 A V-Cycle Multigrid Method

Since the finite element spaces are nested, we can take the coarse-to-fine operator  $I_{k-1}^k: V_{k-1} \longrightarrow V_k$  to be the natural injection. The fine-to-coarse operator  $I_k^{k-1}: V'_k \longrightarrow V'_{k-1}$  is then defined by

$$\langle I_k^{k-1}\ell, \boldsymbol{\nu} \rangle = \langle \ell, I_{k-1}^k \boldsymbol{\nu} \rangle \qquad \forall \ell \in V_k', \, \boldsymbol{\nu} \in V_{k-1}.$$
(7)

We will use a smoother of the form

$$z_{\text{new}} = z_{\text{old}} + M_k^{-1} (g - A_k z_{\text{old}})$$
(8)

for the equation (6), where  $M_k^{-1}: V'_k \longrightarrow V_k$  is a nonoverlapping domain decomposition preconditioner defined below.

#### 3.1 A Nonoverlapping Domain Decomposition Preconditioner

To conform with standard terminology in domain decomposition, in this subsection we will denote  $\mathscr{T}_{k-1}$  by  $\mathscr{T}_{H}$  and  $\mathscr{T}_{k}$  by  $\mathscr{T}_{h}$ . (Thus each element in  $\mathscr{T}_{H}$  is partitioned into eight elements in  $\mathscr{T}_{h}$ ). The spaces  $V_{k-1}$  and  $V_k$  are denoted by  $V_H$  and  $V_h$  respectively. The preconditioner  $M_k^{-1}$  in (8) is denoted by  $M_h^{-1}$  here. It is constructed by substructuring.

For each element  $T \in \mathscr{T}_H$ , we define the twelve dimensional subspace  $V_h^T$  of  $V_h$  by

$$V_h^T = \{ \mathbf{v} \in V_h : \mathbf{v} = 0 \text{ on } \Omega \setminus T \}.$$
(9)

The natural injection from  $V_h^T$  into  $V_h$  is denoted by  $J_T$  and the operator  $A_T : V_h^T \longrightarrow (V_h^T)'$  is defined by

$$\langle A_T \boldsymbol{w}, \boldsymbol{v} \rangle = a(\boldsymbol{w}, \boldsymbol{v}) \qquad \forall \boldsymbol{v}, \boldsymbol{w} \in V_h^T.$$
 (10)

Let  $\mathscr{F}_H$  be the set of the interior faces of the triangulation  $\mathscr{T}_H$ . Given any  $F \in \mathscr{F}_H$  that is the common face of two elements  $T_F^+$  and  $T_F^-$  in  $\mathscr{T}_H$ , we define the four dimensional subspace  $V_h^F$  of  $V_h$  by

$$V_h^F = \{ \boldsymbol{\nu} \in V_h : \boldsymbol{\nu} = \boldsymbol{0} \text{ on } \boldsymbol{\Omega} \setminus (T_F^- \cup T_F^+) \text{ and } a(\boldsymbol{\nu}, \boldsymbol{w}) = 0 \quad \forall \, \boldsymbol{w} \in (V_h^{T_F^-} + V_h^{T_F^+}) \}.$$
(11)

The natural injection from  $V_h^F$  into  $V_h$  is denoted by  $J_F$  and the operator  $A_F : V_h^F \longrightarrow (V_h^F)'$  is defined by

$$\langle A_F \boldsymbol{w}, \boldsymbol{v} \rangle = a(\boldsymbol{w}, \boldsymbol{v}) \qquad \forall \boldsymbol{v}, \boldsymbol{w} \in V_h^F.$$
 (12)

if  $\boldsymbol{w} \in V_h$  has the same degrees of freedom as  $\boldsymbol{v}$  on  $\partial T_F^+ \cup \partial T_F^-$ .

The subspaces associated with the elements and interior faces of  $\mathcal{T}_H$  form a direct sum decomposition of  $V_h$ :

$$V_h = \sum_{T \in \mathscr{T}_H} V_h^T + \sum_{F \in \mathscr{F}_H} V_h^F,$$
(13)

and the preconditioner  $M_h^{-1}$  is given by

$$M_h^{-1} = \eta_F \left(\sum_{T \in \mathscr{T}_H} J_T A_T^{-1} J_T^t + \sum_{F \in \mathscr{T}_H} J_F A_F^{-1} J_F^t\right),\tag{14}$$

where  $\eta_F$  is a damping factor and  $J_T^t : V_h' \longrightarrow (V_h^T)'$  (resp.  $J_F^t : V_h' \longrightarrow (V_h^F)'$ ) is the transpose of  $J_T$  (resp.  $J_F$ ) with respect to the canonical bilinear forms.

## 3.2 The k<sup>th</sup> Level V-Cycle Multigrid Algorithm

The output  $MG(k, g, z_0, m)$  of the  $k^{\text{th}}$  level (symmetric) multigrid *V*-cycle algorithm for (6), with initial guess  $z_0 \in V_k$  and *m* smoothing steps, is defined by the following recursive steps:

For k = 0, the output is obtained from a direct method:

$$MG(0,g,\mathbf{z}_0,m) = A_0^{-1}g.$$

For  $k \ge 1$ , we set

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$$\begin{aligned} \mathbf{z}_{l} &= \mathbf{z}_{l-1} + M_{k}^{-1} \left( g - A_{k} \mathbf{z}_{l-1} \right) & \text{for } 1 \leq l \leq m, \\ \overline{g} &= I_{k}^{k-1} \left( g - A_{k} \mathbf{z}_{m} \right), \\ \mathbf{z}_{m+1} &= \mathbf{z}_{m} + I_{k-1}^{k} MG \left( k - 1, \overline{g}, 0, m \right), \\ \mathbf{z}_{l} &= \mathbf{z}_{l-1} + M_{k}^{-1} \left( g - A_{k} \mathbf{z}_{l-1} \right) & \text{for } m+2 \leq l \leq 2m+1. \end{aligned}$$

The output of  $MG(k, g, z_0, m)$  is  $z_{2m+1}$ .

*Remark 1.* Given  $\ell \in V'_k$ , the cost of computing  $M_k^{-1}\ell$  is  $O(n_k)$ , where  $n_k$  is the dimension of  $V_k$ . Therefore the overall cost for computing  $MG(k, g, \mathbf{z}_0, m)$  is also  $O(n_k)$ .

If the domain  $\Omega$  is convex, we have the following convergence theorem:

**Theorem 1.** If  $z \in V_k$  and  $g \in V'_k$  satisfy  $A_k z = g$ , then we have

$$\|\mathbf{z} - MG(k, g, \mathbf{z}_0, m)\|_a \leq \frac{C}{C + 2m} \|\mathbf{z} - \mathbf{z}_0\|_a \qquad \forall k \geq 1,$$

*where*  $\| \cdot \|_{a}^{2} = a(\cdot, \cdot).$ 

Due to space restriction, a detailed analysis will not be reported here. Further details are provided in [5].

## **4** Numerical Results

### 4.1 Jump Coefficient



Fig. 1: Checkerboard distribution of the coefficients

In the first experiment we consider (1) on the unit cube  $\Omega = (0,1)^3$ . We apply multigrid algorithms with smoothers introduced in Section 3.1. The damping factor  $\eta_F$  is taken to be 1/11. The initial triangulation  $\mathscr{T}_0$  consists of eight identical cubes and we use the coefficients  $\alpha$  and  $\beta$  that have jumps across the interface between the sub-cubes with a checkerboard pattern as in Fig. 1. We estimate the contraction numbers of the  $k^{\text{th}}$  level V-cycle multigrid method for k = 1, ..., 5 and for m smoothing steps, where m = 1, ..., 6. We report the contraction numbers obtained by computing the largest eigenvalue of the error propagation operators. The results are presented in Table 1. The uniform convergence of the V-cycle multigrid methods for  $m \ge 1$  is clearly observed and the method is not sensitive to the jumps of coefficients.

Table 1: Contraction numbers of the V-cycle multigrid method for the unit cube.  $\alpha_b$  and  $\beta_b$  for the black subregions and  $\alpha_w$  and  $\beta_w$  for the white subregions as indicated in a checkerboard pattern as

	e
	m = 1 $m = 2$ $m = 3$ $m = 4$ $m = 5$ $m = 6$
	$\alpha_b = 0.01, \beta_b = 100, \alpha_w = 1, \beta_w = 1$
k = 1	8.3e-1 6.8e-1 4.7e-1 2.2e-1 5.1e-2 4.7e-3
k = 2	9.0e-1 8.2e-1 7.1e-1 5.1e-1 3.2e-1 2.7e-1
k = 3	9.3e-1 8.8e-1 7.9e-1 6.4e-1 5.2e-1 4.7e-1
k = 4	9.3e-1 9.0e-1 8.4e-1 7.2e-1 6.4e-1 6.0e-1
k = 5	9.3e-1 9.0e-1 8.6e-1 7.8e-1 6.9e-1 6.9e-1
	$\alpha_b = 0.1, \beta_b = 10, \alpha_w = 1, \beta_w = 1$
k = 1	8.7e-1 7.7e-1 6.0e-1 3.8e-1 2.1e-1 8.1e-2
k = 2	9.1e-1 8.4e-1 7.1e-1 5.4e-1 3.6e-1 2.8e-1
k = 3	9.2e-1 8.7e-1 7.8e-1 6.4e-1 5.2e-1 4.7e-1
k = 4	9.3e-1 9.0e-1 8.4e-1 7.4e-1 6.5e-1 6.0e-1
k = 5	9.4e-1 9.1e-1 8.7e-1 8.0e-1 7.2e-1 6.9e-1
	$\alpha_b = 1, \beta_b = 1, \alpha_w = 1, \beta_w = 1$
k = 1	9.1e-1 8.3e-1 7.1e-1 5.0e-1 3.1e-1 2.3e-1
k = 2	9.2e-1 8.7e-1 7.9e-1 6.3e-1 5.0e-1 4.3e-1
k = 3	9.3e-1 9.0e-1 8.4e-1 7.4e-1 6.3e-1 5.8e-1
k = 4	9.4e-1 9.1e-1 8.7e-1 8.0e-1 7.1e-1 6.7e-1
k = 5	9.4e-1 9.2e-1 8.8e-1 8.2e-1 7.5e-1 7.2e-1
	$\alpha_b = 10, \beta_b = 0.1, \alpha_w = 1, \beta_w = 1$
k = 1	9.0e-1 8.4e-1 7.0e-1 4.9e-1 3.3e-1 2.8e-1
k = 2	9.2e-1 8.9e-1 7.9e-1 6.4e-1 5.2e-1 4.7e-1
k = 3	9.4e-1 9.1e-1 8.4e-1 7.4e-1 6.4e-1 6.0e-1
k = 4	9.4e-1 9.1e-1 8.6e-1 8.0e-1 7.3e-1 6.8e-1
k = 5	9.4e-1 9.2e-1 8.9e-1 8.2e-1 7.6e-1 7.4e-1
	$\alpha_b = 100, \beta_b = 0.01, \alpha_w = 1, \beta_w = 1$
k = 1	9.1e-1 8.4e-1 7.1e-1 5.1e-1 3.3e-1 2.9e-1
$k = \overline{2}$	9.3e-1 8.9e-1 7.9e-1 6.5e-1 5.2e-1 4.8e-1
$k = \overline{3}$	9.3e-1 9.1e-1 8.5e-1 7.4e-1 6.4e-1 6.0e-1
k = 4	9.4e-1 9.2e-1 8.8e-1 8.0e-1 7.1e-1 6.9e-1
$k = \overline{5}$	9.4e-1 9.3e-1 9.0e-1 8.4e-1 7.7e-1 7.5e-1

in Fig. 1

### 4.2 Nonconvex Domain



Fig. 2: Nonconvex domain

In the second numerical experiment we report the results for our model problem (1) on the nonconvex domain  $\Omega = (0,1)^3 \setminus ([1/2,1]^3)$ . We use the constant coefficients  $\alpha = 1$  and  $\beta = 1$  and other general settings are quite similar to those of Section 4.1. The results are presented in Table 2. It is observed that the method provides a uniform convergence of the *V* cycle multigrid. However, the contraction numbers are generally larger that those of the convex domain.

Table 2: Contraction numbers of the V-cycle multigrid method for the non-convex domain as in Fig. 2 with  $\alpha = 1, \beta = 1$ 

	m = 1	m=2	m = 3	m = 4	m = 5	m = 6
k = 1	9.3e-1	9.0e-1	8.2e-1	6.9e-1	5.5e-1	4.6e-1
k = 2	9.5e-1	9.2e-1	8.5e-1	7.7e-1	6.8e-1	6.3e-1
<i>k</i> = 3	9.6e-1	9.2e-1	8.8e-1	8.2e-1	7.7e-1	7.3e-1
k = 4	9.6e-1	9.3e-1	8.9e-1	8.5e-1	8.0e-1	7.8e-1
k = 5	9.6e-1	9.3e-1	9.0e-1	8.7e-1	8.4e-1	8.2e-1

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